

Particle acceleration in solar flares: merging magnetic islands in forced reconnection

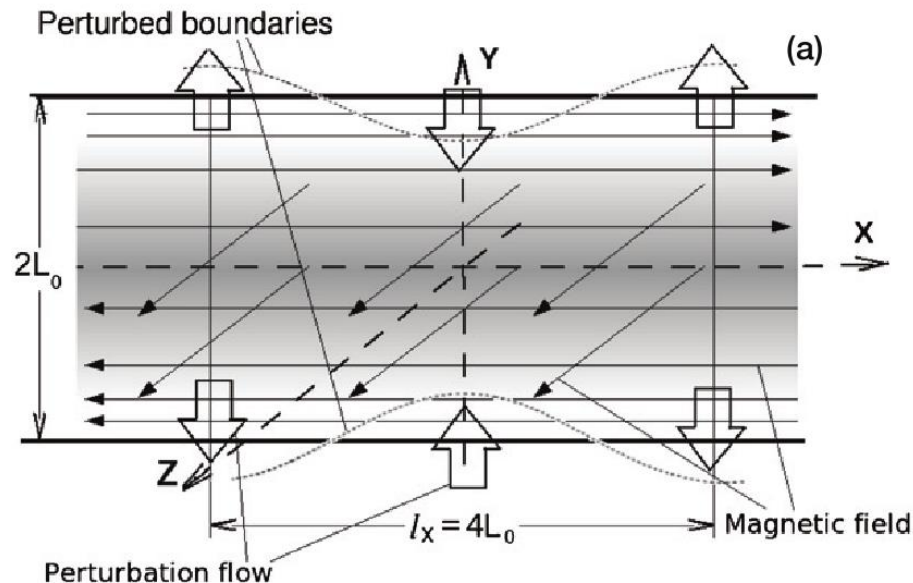
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Image: Anthony Holloway

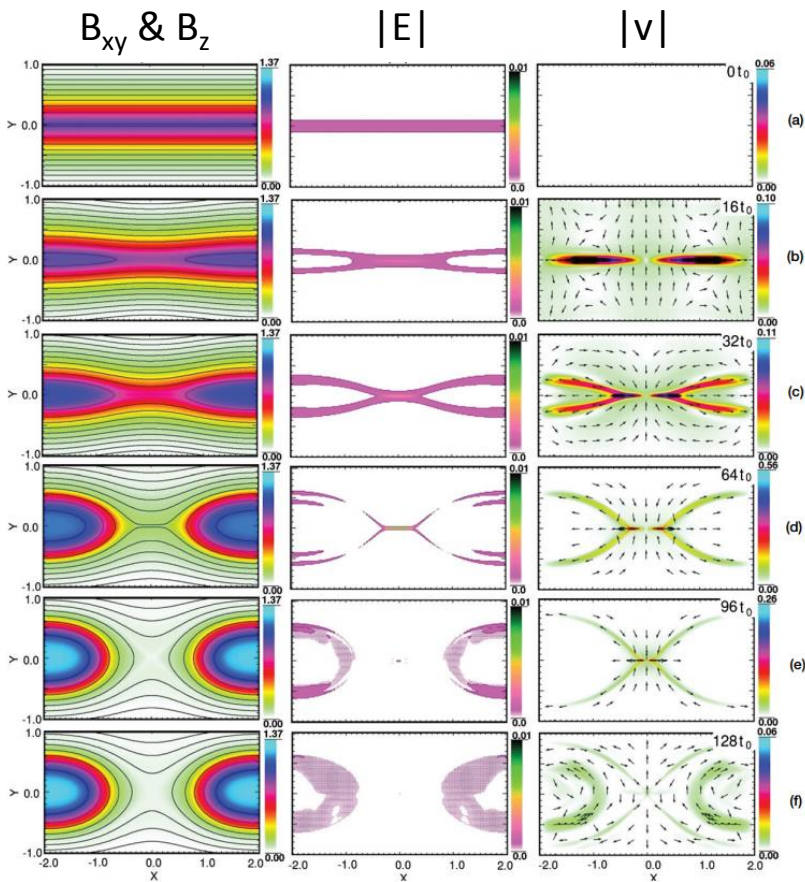
Forced magnetic reconnection

- 2D slab of plasma with sheared force-free field stable to tearing mode
 - *Hahm & Kulsrud 1985, Vekstein & Jain 1998*
- Apply transient perturbation to the boundaries
 - Analytic models: directly perturb the boundary of the plasma
 - Numerical models: normal flows at boundary
- Magnetic reconnection forms a chain of magnetic islands

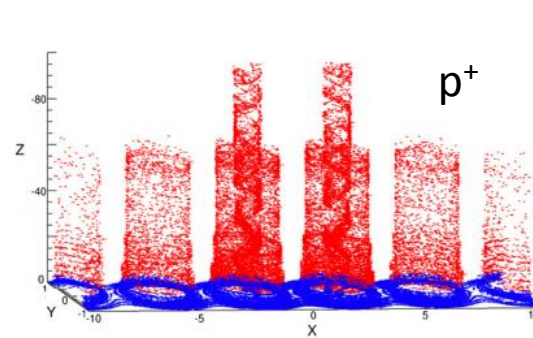


Forced reconnection diagram – *Gordovskyy et al 2010*

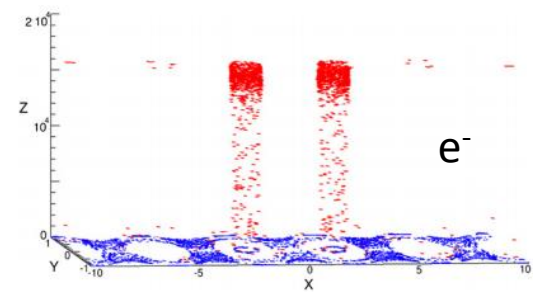
Particle acceleration in forced reconnection: 2D MHD + test particles



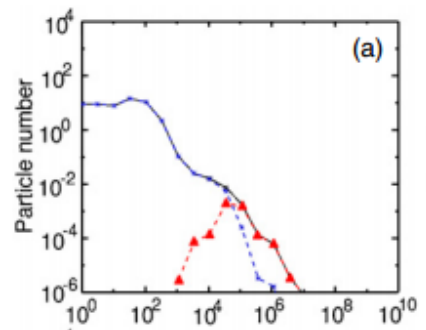
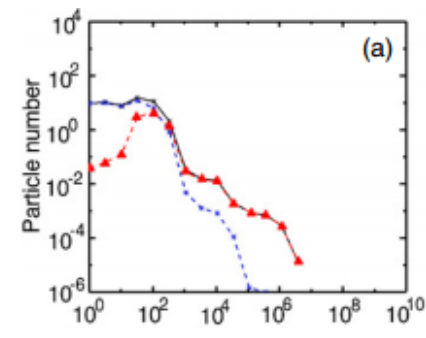
MHD simulation for forced reconnection



Test particle final positions



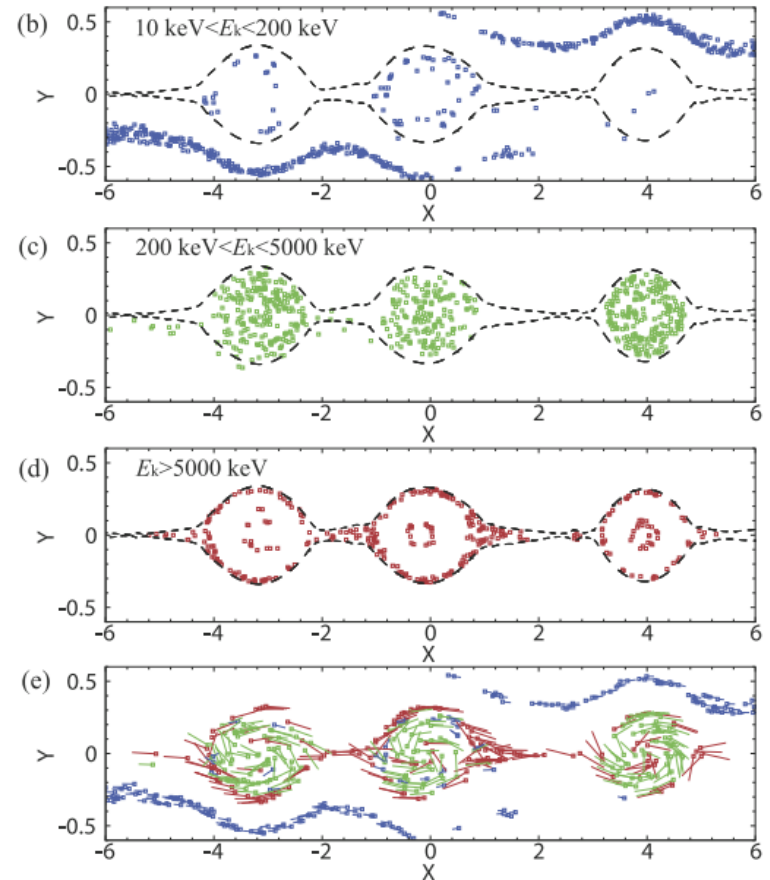
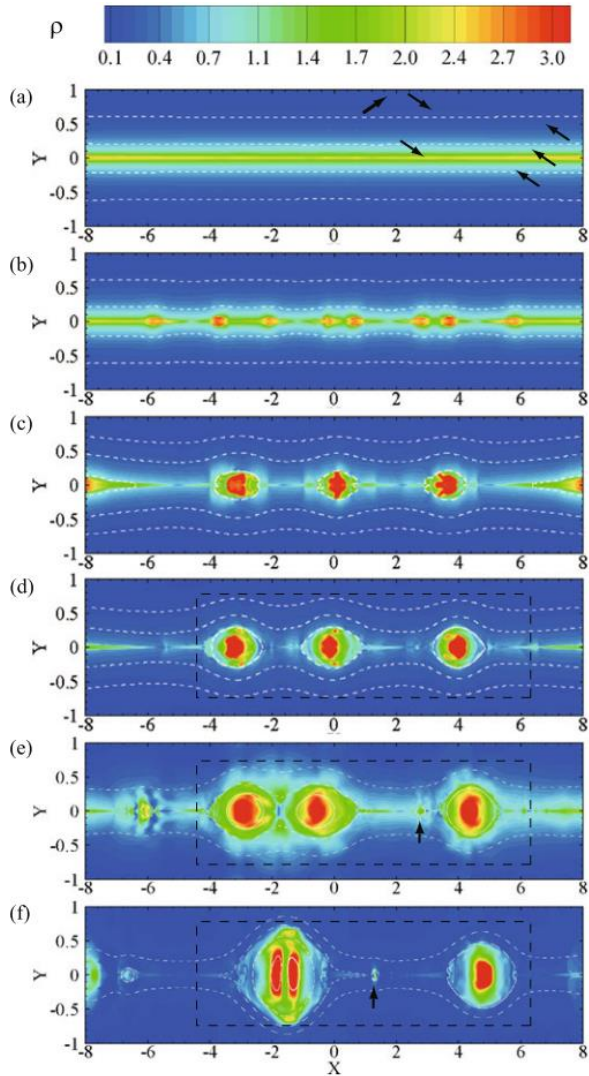
Test particle final positions



Particle spectra

Merging islands in tearing unstable current sheet

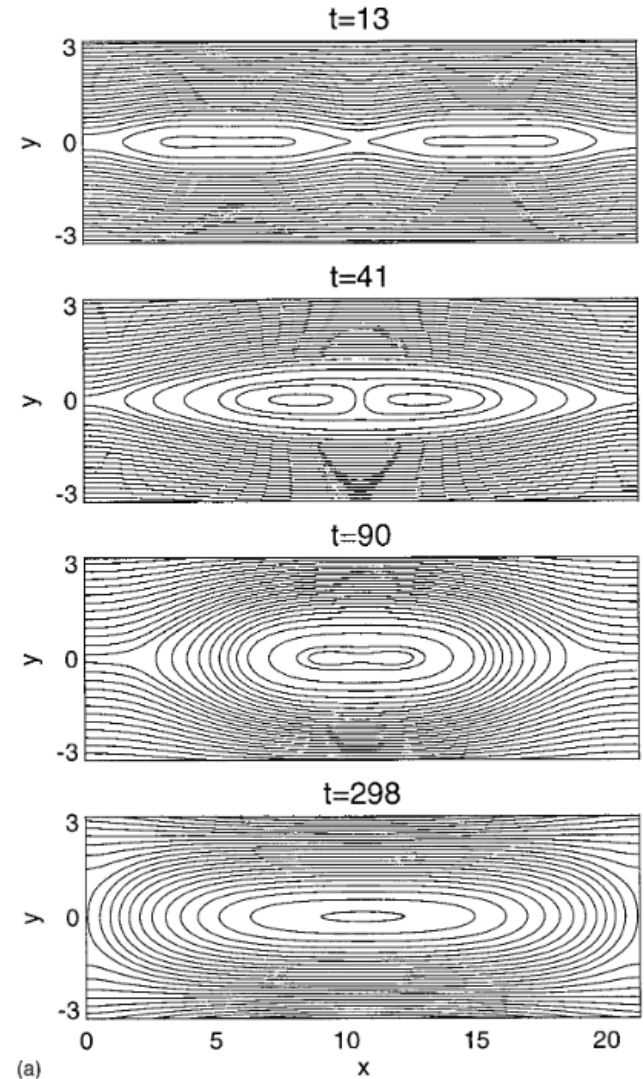
MHD + Test Particle approach



Test particles in stationary field snapshots
 – Zhang 2015

Coalescence

- Chain of magnetic islands (O-points)
- Attractive parallel currents
- Neighbouring islands coalesce
- To simulate this:
 - Capture multiple islands and allow symmetry-breaking with a long numerical box
 - Simulate large period of time (> 150 Alfvén times)



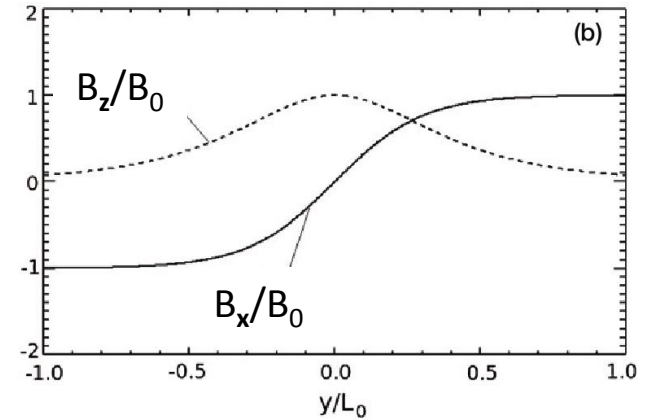
Aims of current study

- To investigate effect of different spatial driving disturbances on forced reconnection – field evolution and energetic particles
 - In general, driving disturbance is superposition of many wavelengths
 - How do islands develop and evolve for multi-wavelength perturbations?
 - Does forced reconnection “work” for more realistic driving disturbances e.g. localised perturbation?
- To investigate reconnection, energy release and particle acceleration during island formation and island coalescence
 - How does distribution of energetic particles evolve through different phases?
 - How does merger of islands affect particle energy spectra and spatial distributions?

Simulation set up

- Initial 1D force-free field:

$$\mathbf{B}_i = \left\{ B_0 \tanh \left(\frac{y}{y_0} \right), 0, B_0 \cosh^{-1} \left(\frac{y}{y_0} \right) \right\}$$



- Boundary conditions:

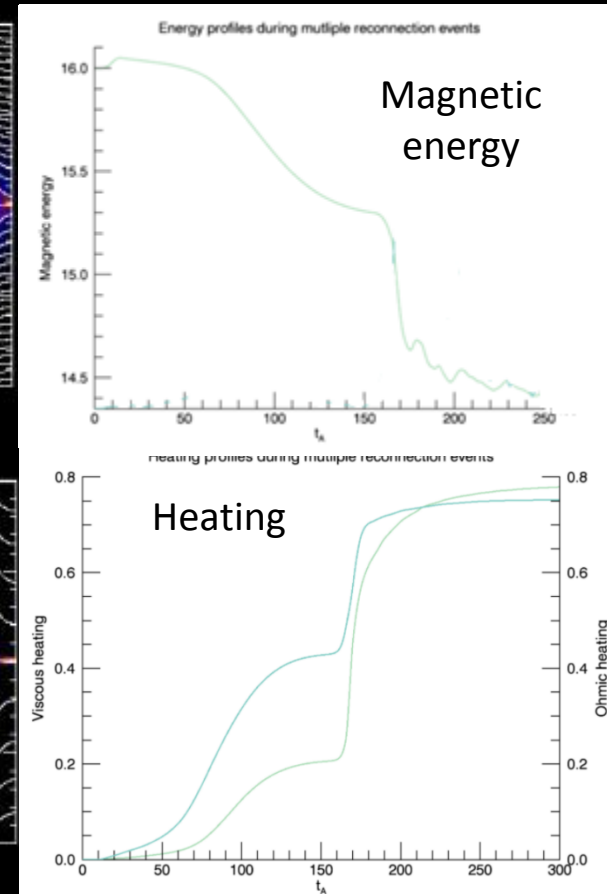
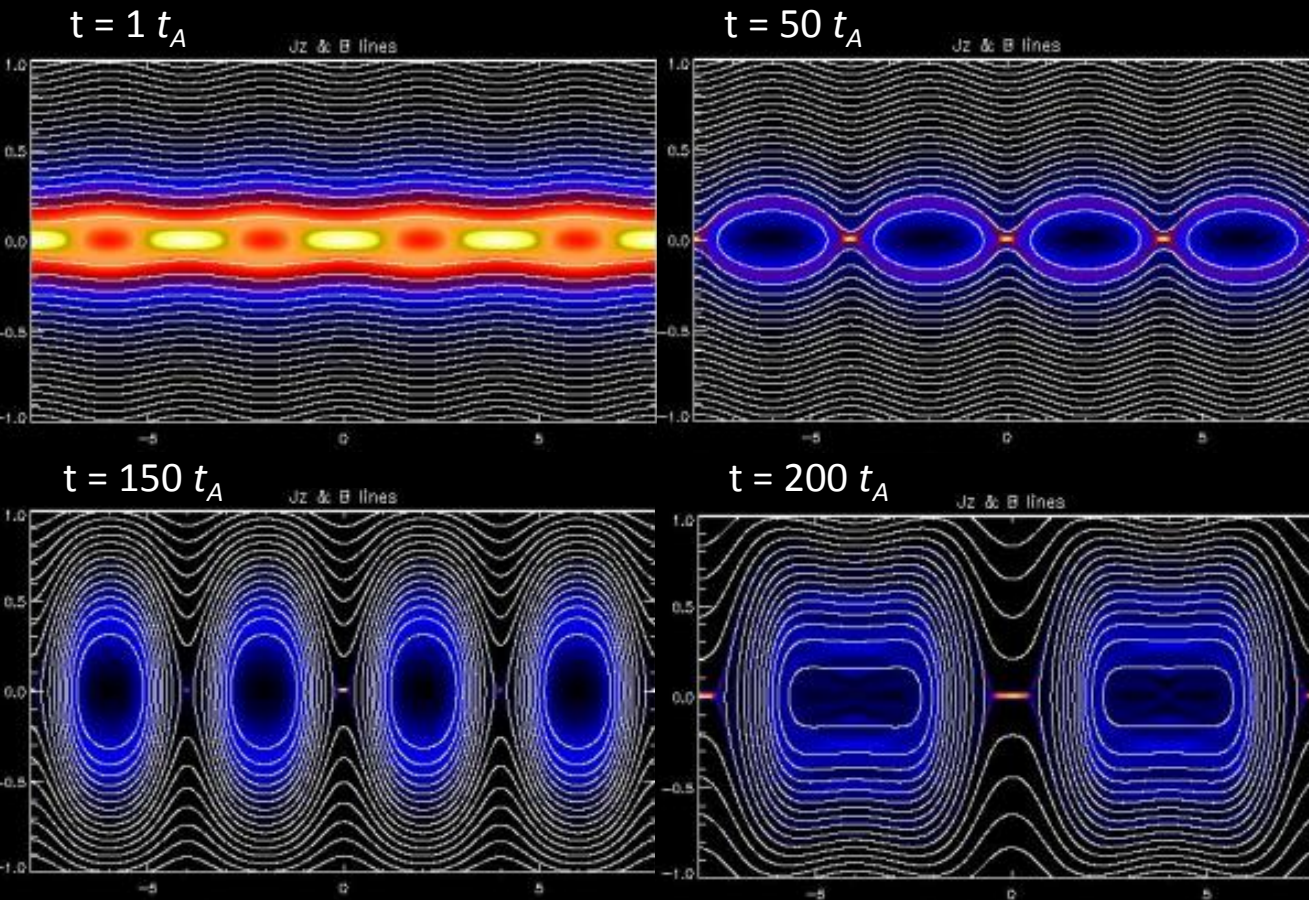
- Conducting walls in y $y = \pm a$: $v = B_z = \frac{\partial B_x}{\partial y} = 0$, $\frac{\partial B_x}{\partial y} = \frac{\partial B_y}{\partial x}$
- Periodic in x

- Transient spatially-sinusoidal boundary perturbation:

$$t < \tau_p : \quad v_y = \frac{\Delta}{\tau_p} \left[1 - \cos \left(\frac{2\pi t}{\tau_p} \right) \right] \cos \left(\frac{2\pi x}{L} \right)$$

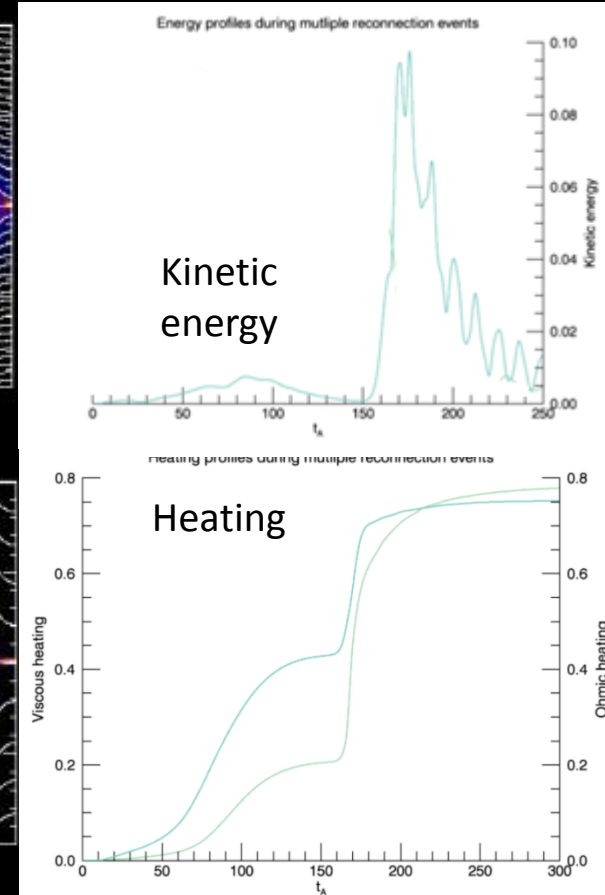
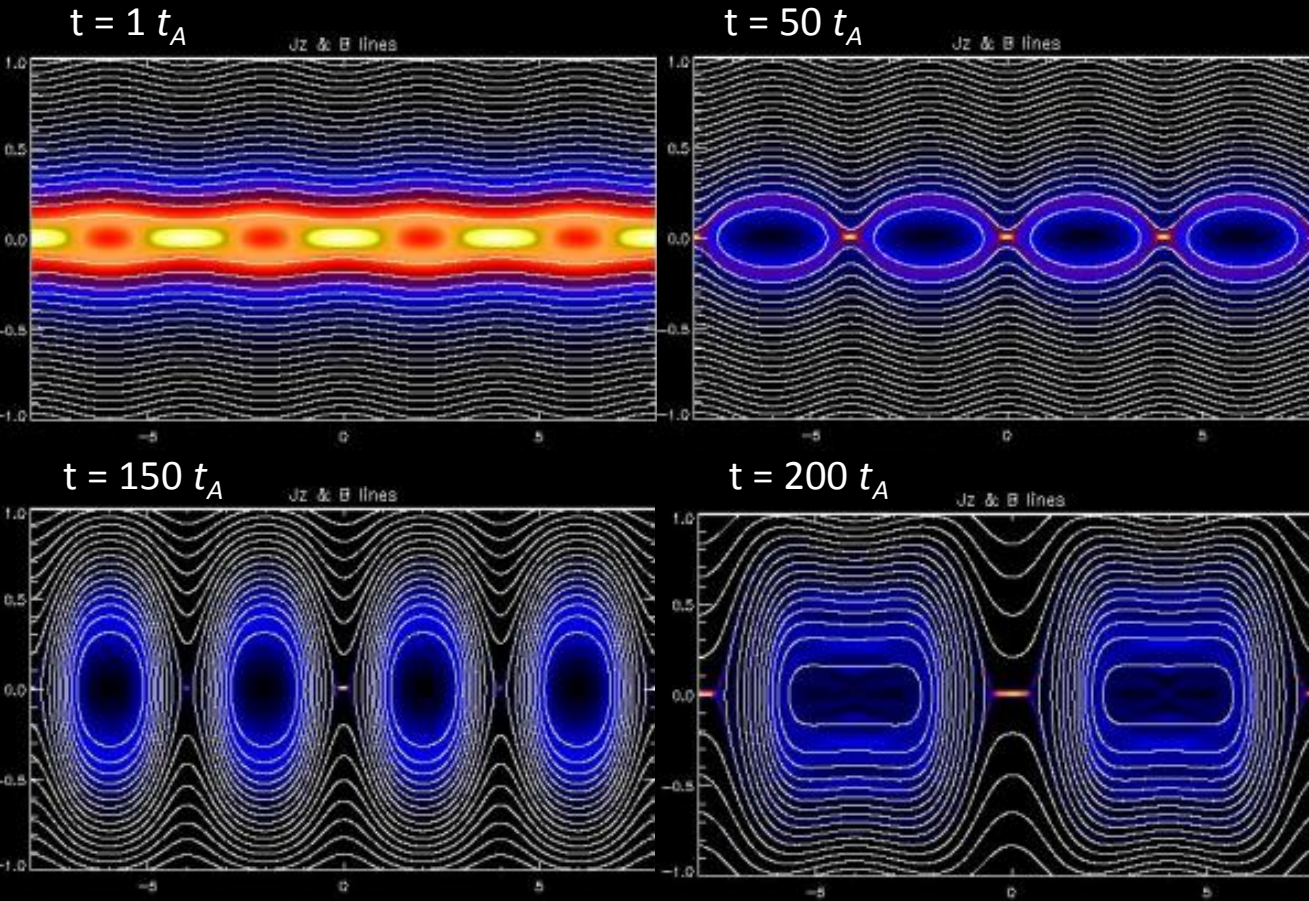
- Simulations were performed using **Lare2D** (Arber et al. 2001)
 - **Anomalous resistivity** was used: resistivity is enhanced in regions of strong current.

Long cell simulations



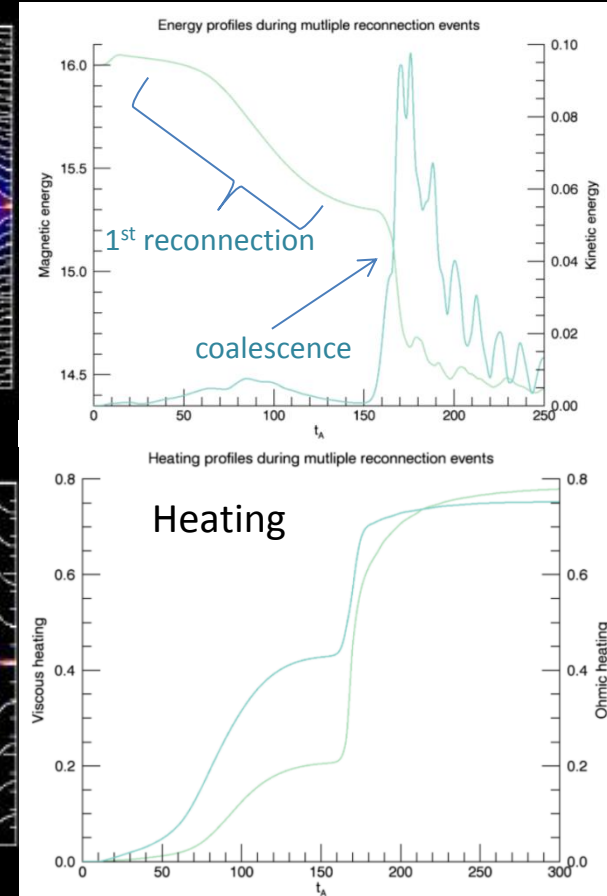
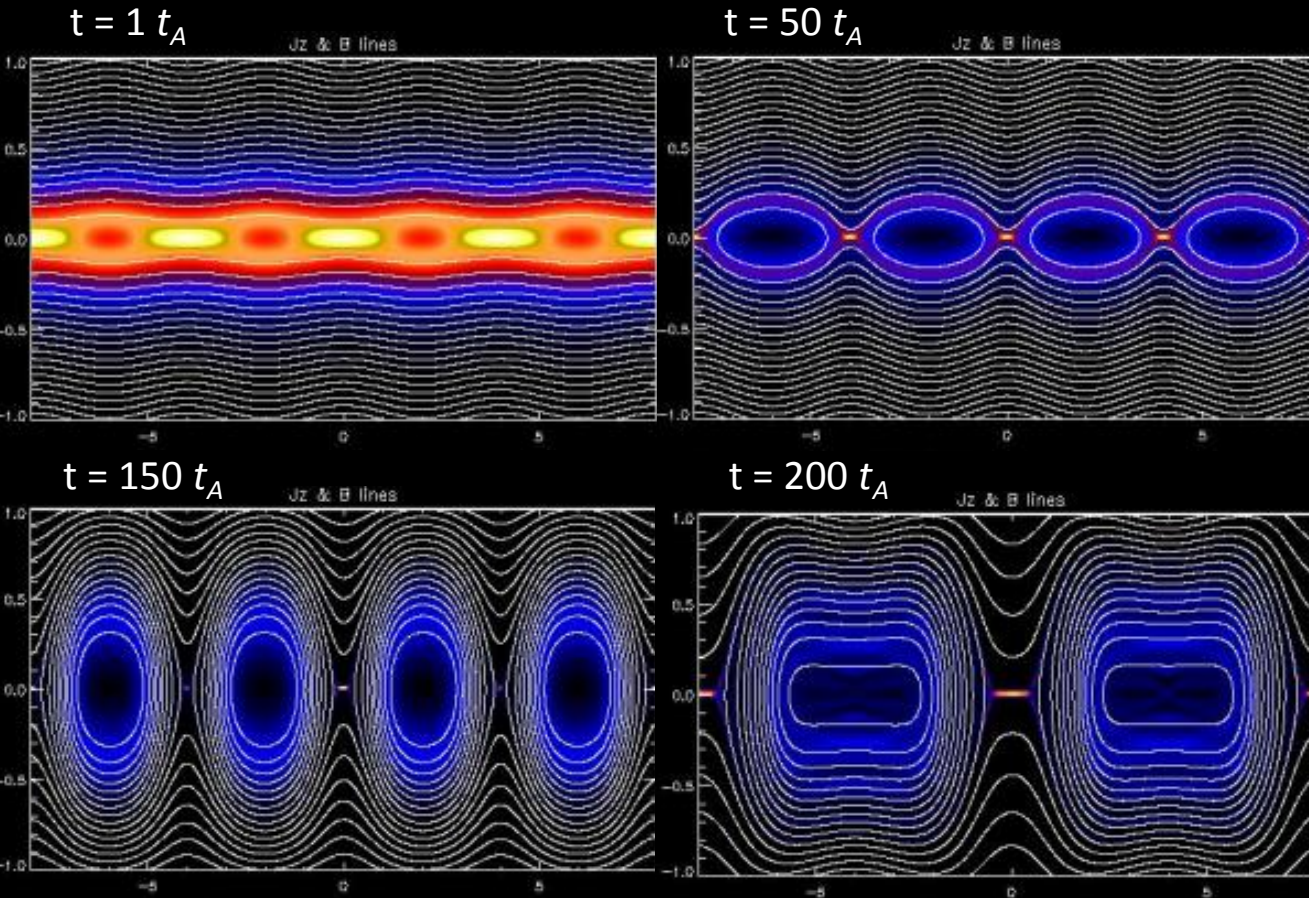
- 16 by 2 length grid with 2048 by 256 grid points
- **Resistivity:** $\eta = 10^{-4}$ when $J > 6$, else $\eta = 10^{-7}$
- **Perturbation:** amplitude $\Delta = 0.05$, wavelength $L = 4$, duration $t_p = 16 t_A$

Long cell simulations



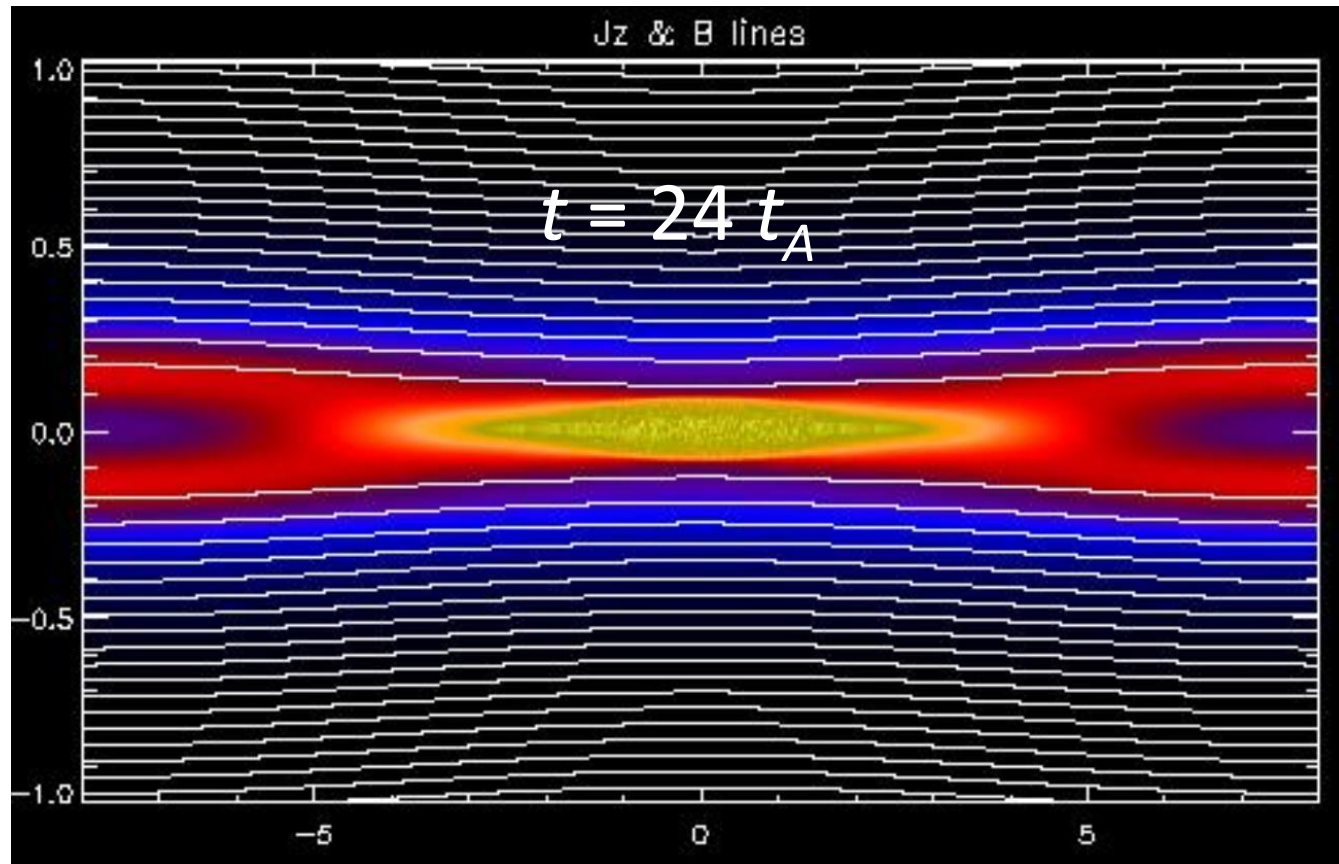
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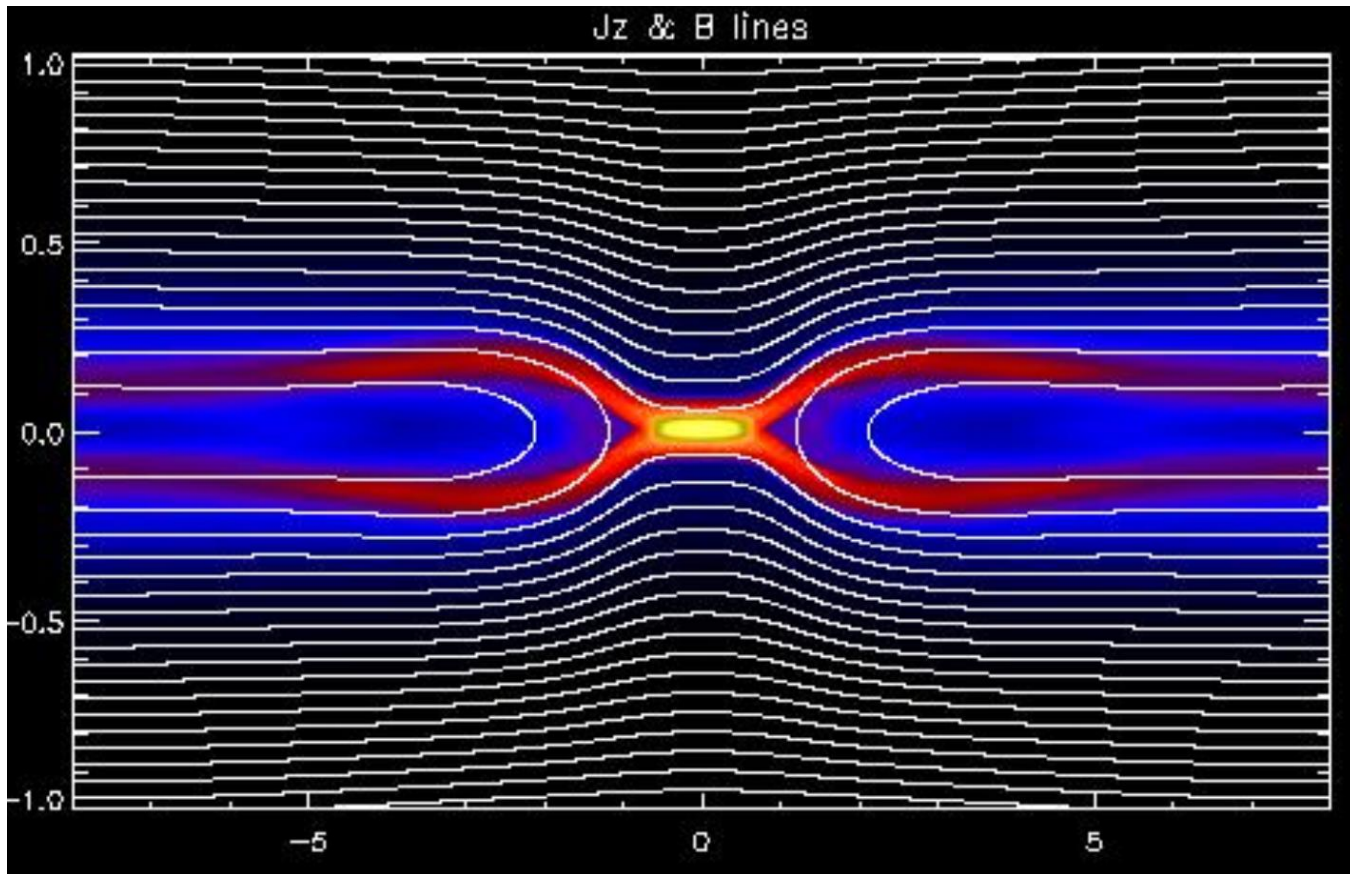


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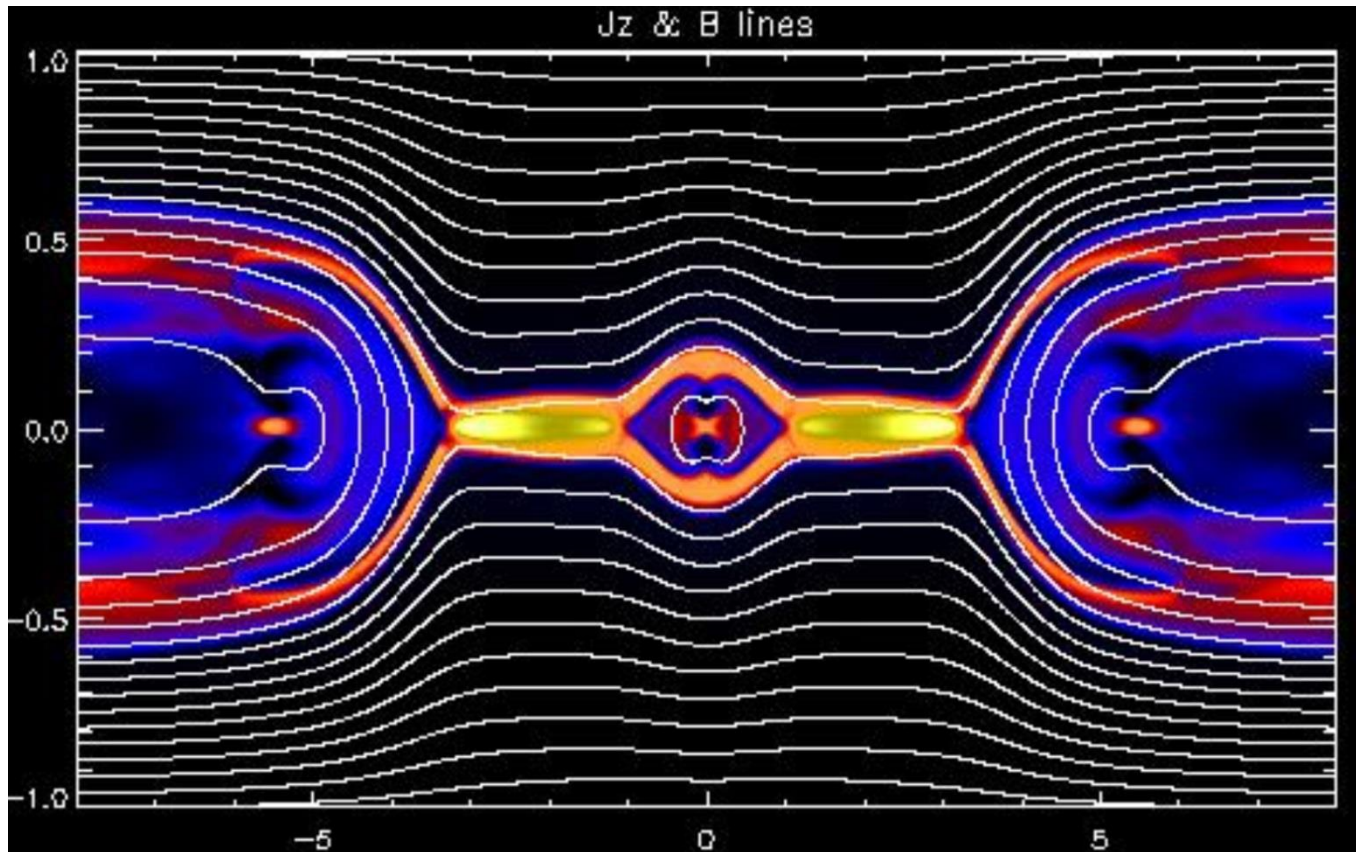
Single long-wavelength driving disturbance



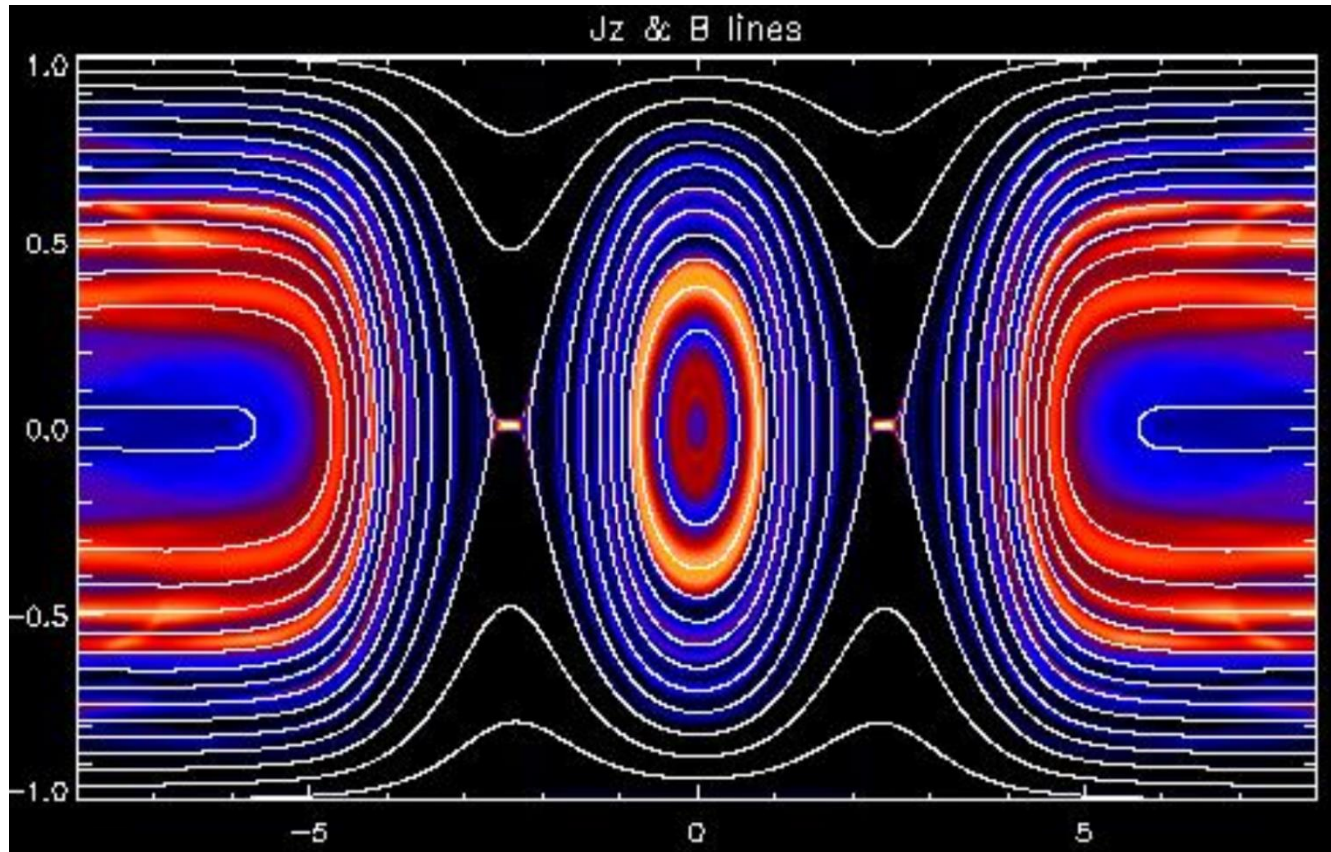
Single long-wavelength driving disturbance



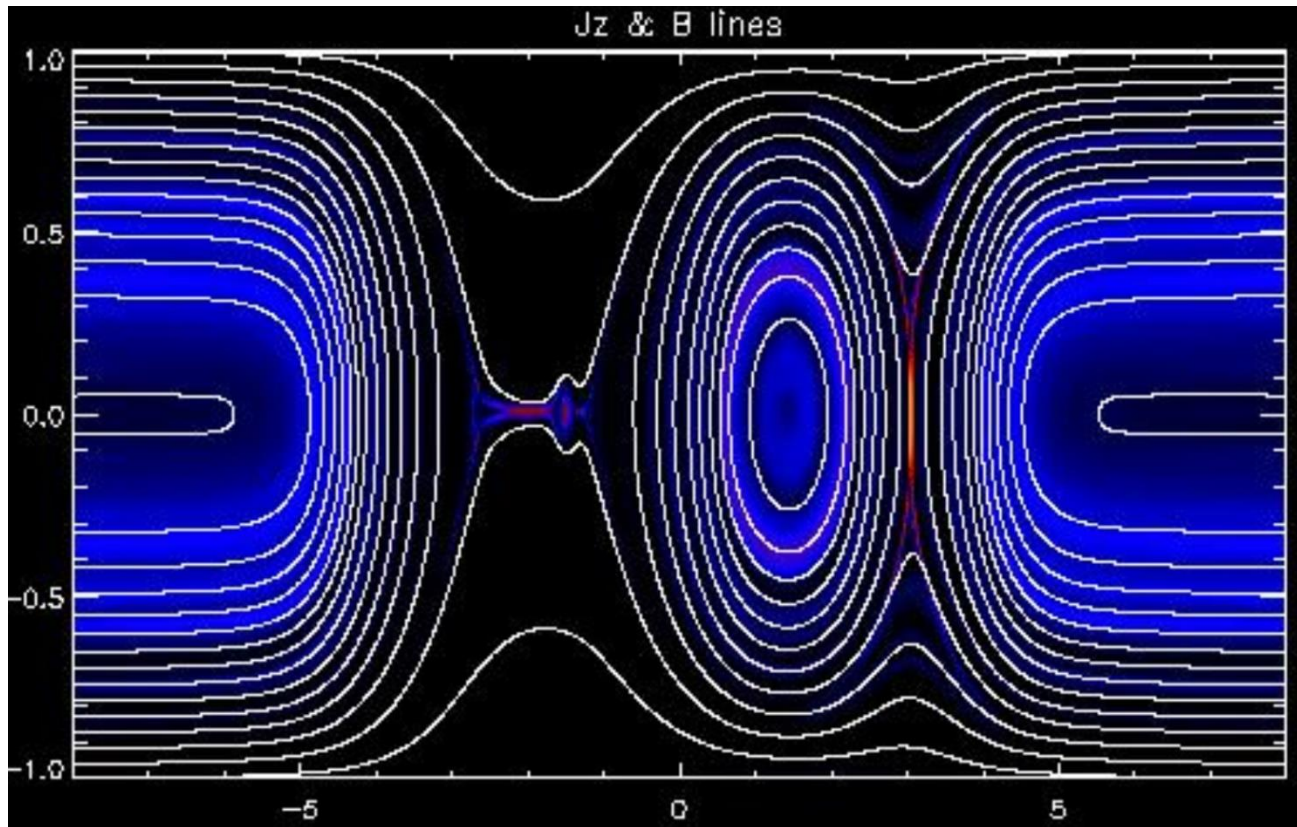
Single long-wavelength driving disturbance



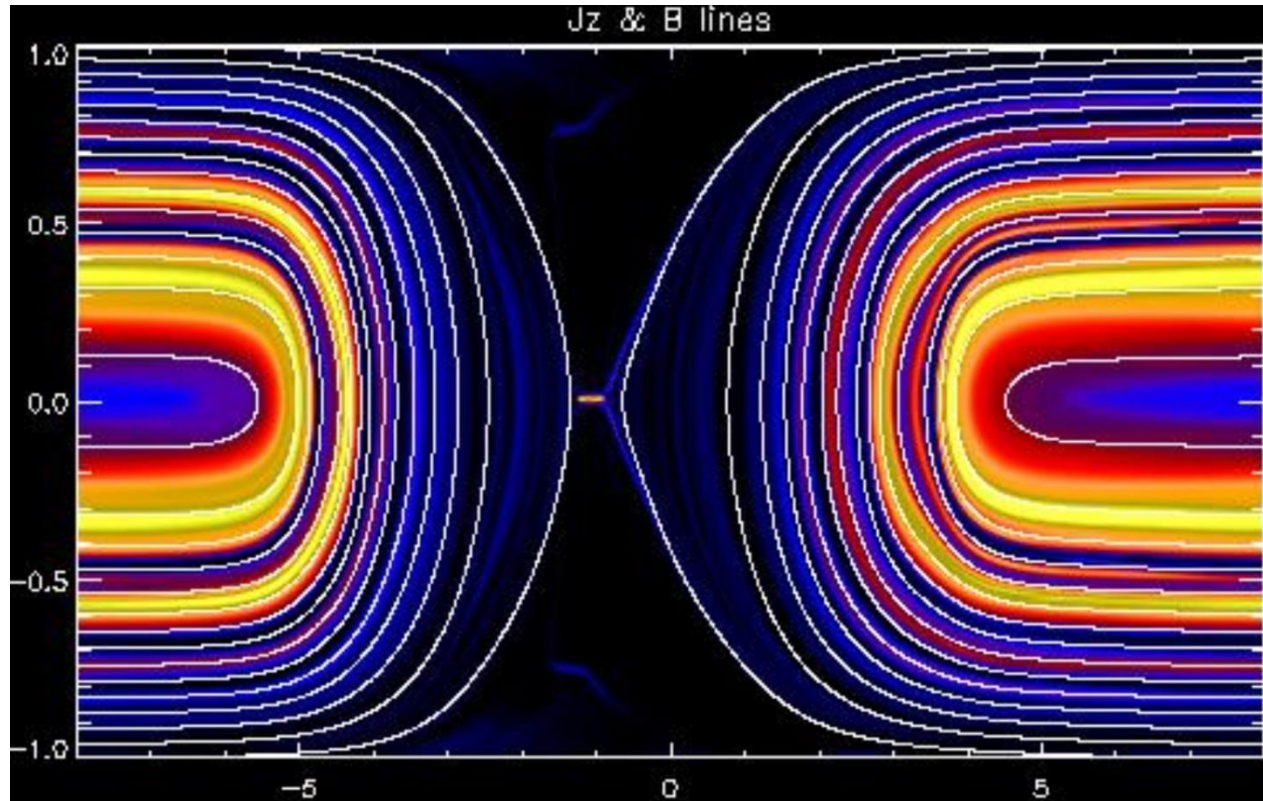
Single long-wavelength driving disturbance



Single long-wavelength driving disturbance

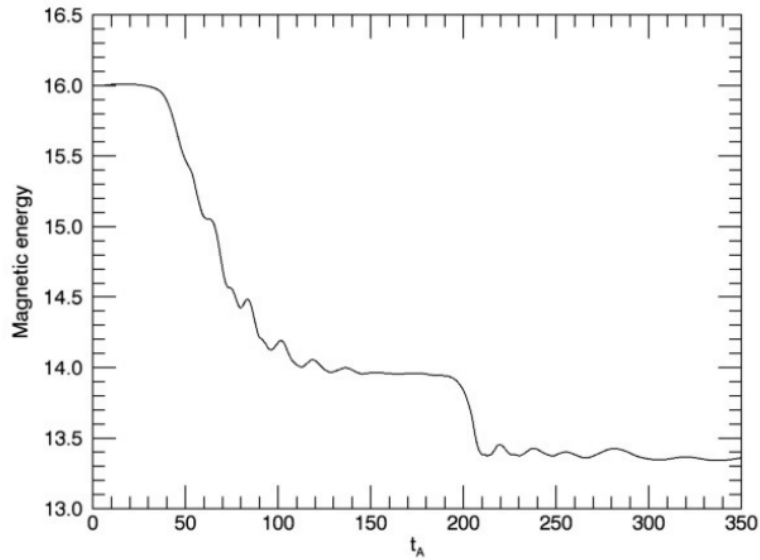


Single long-wavelength driving disturbance

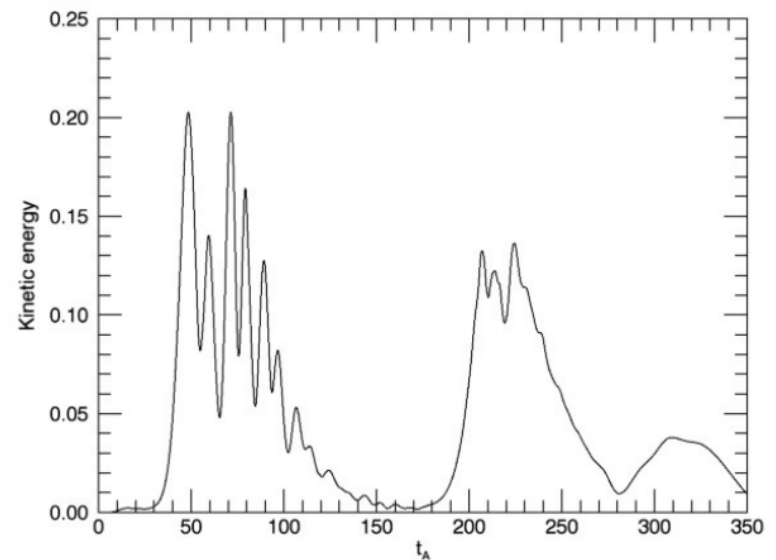


Energetics with merging islands

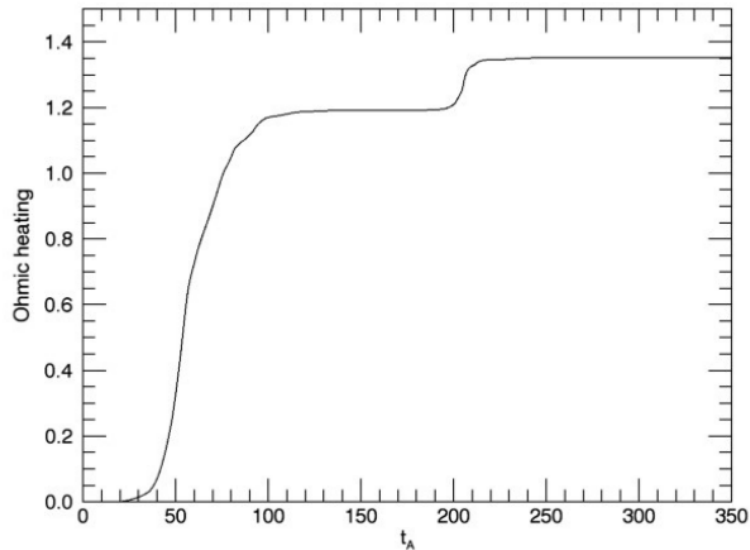
Magnetic energy



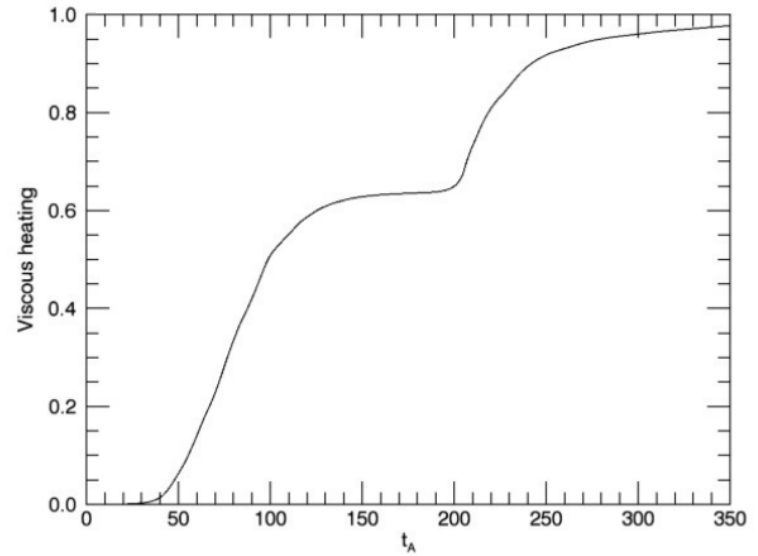
Kinetic energy



Ohmic heating

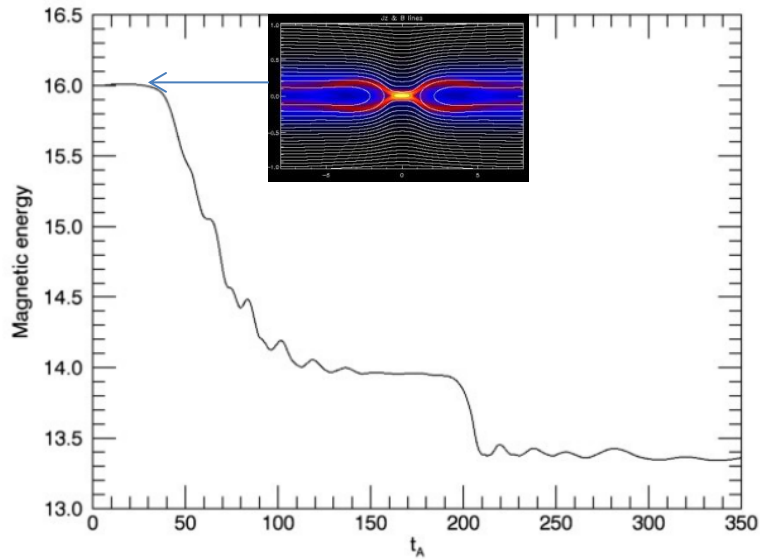


Viscous heating

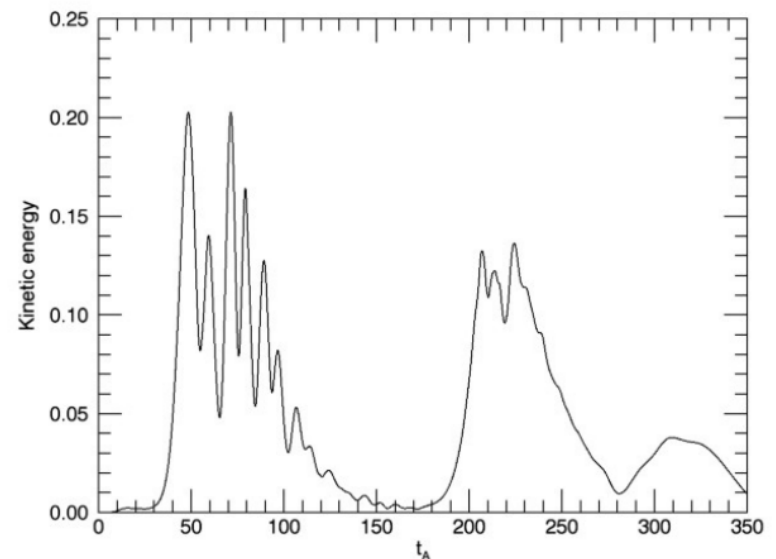


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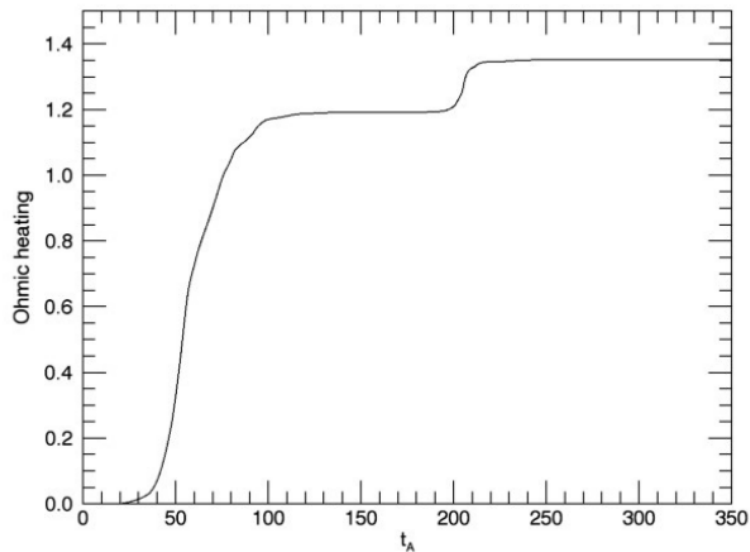
Magnetic energy



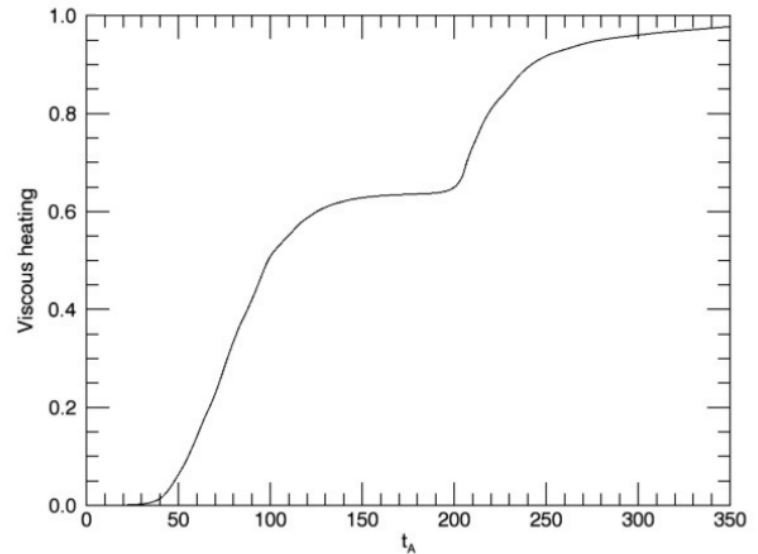
Kinetic energy



Ohmic heating

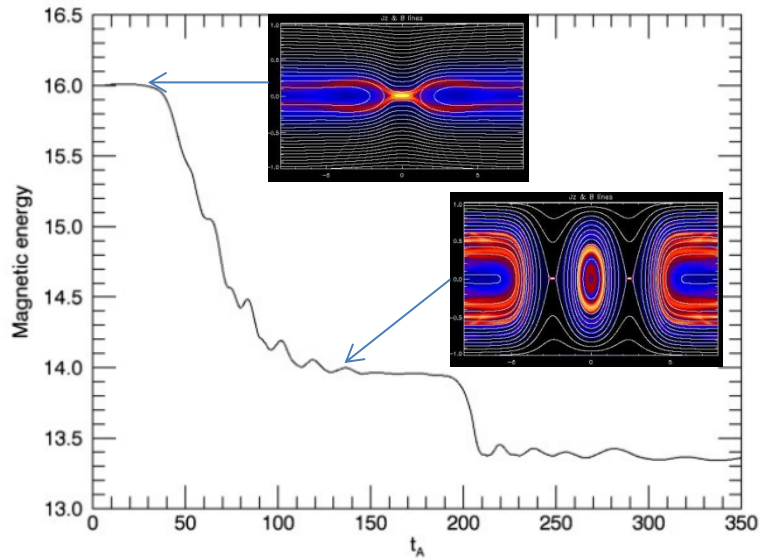


Viscous heating

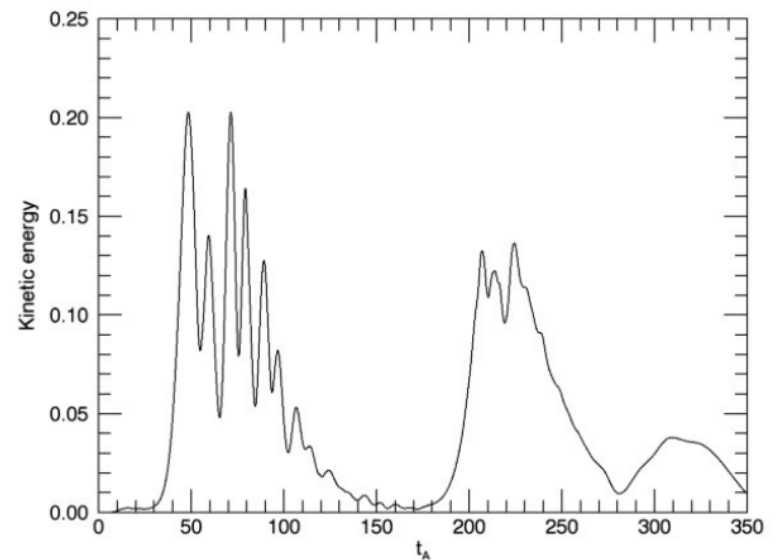


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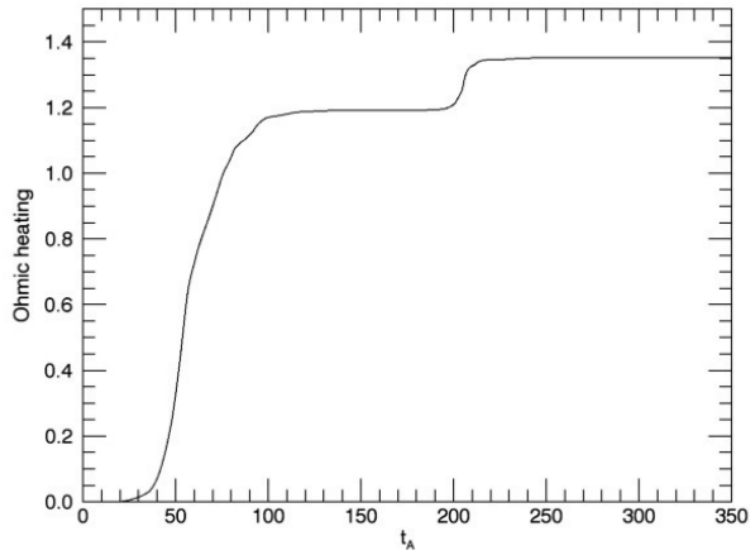
Magnetic energy



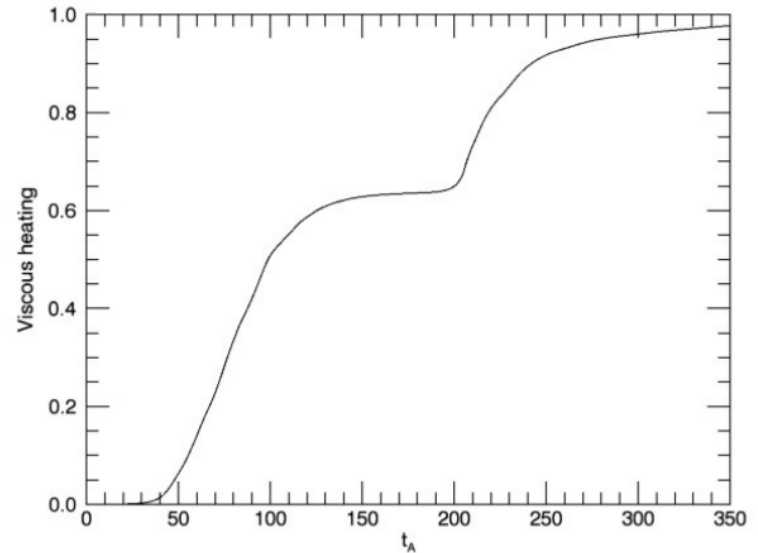
Kinetic energy



Ohmic heating

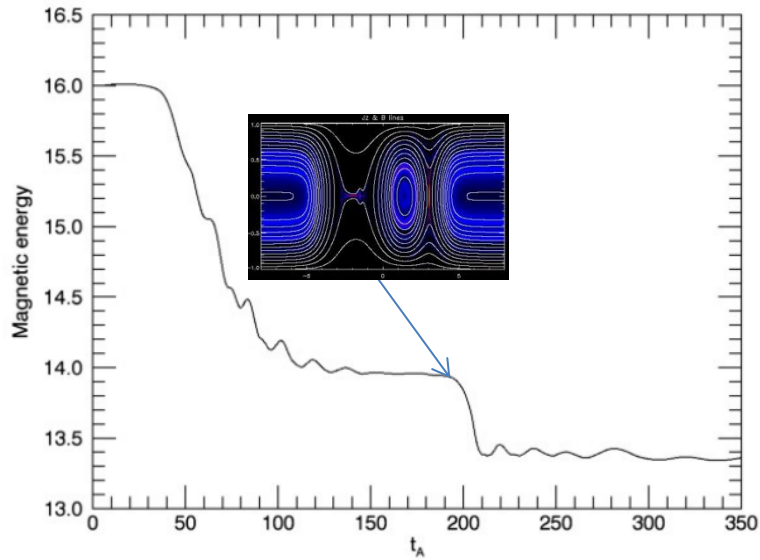


Viscous heating

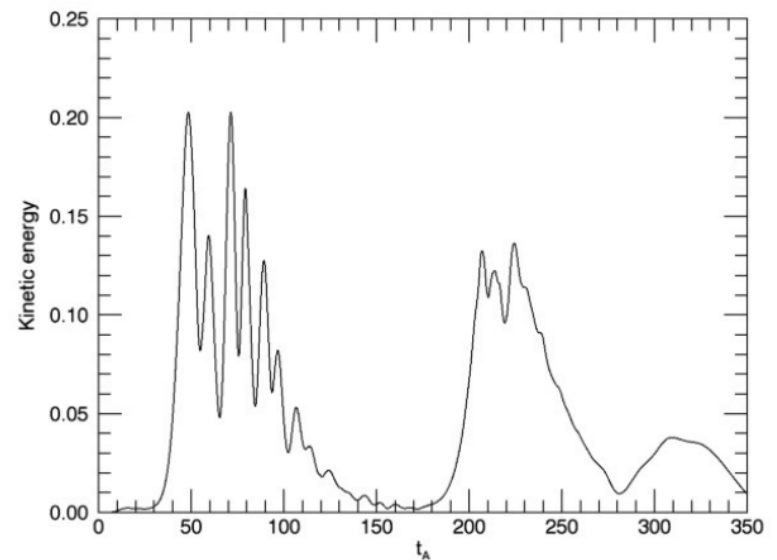


Energetics with merging islands

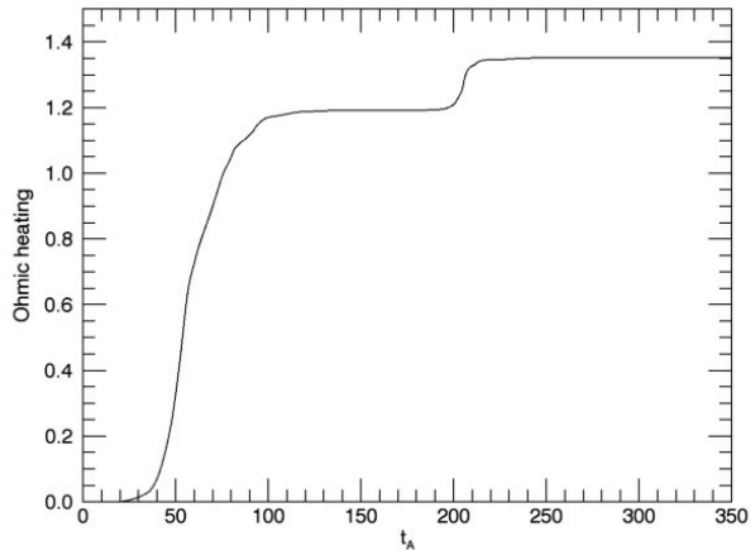
Magnetic energy



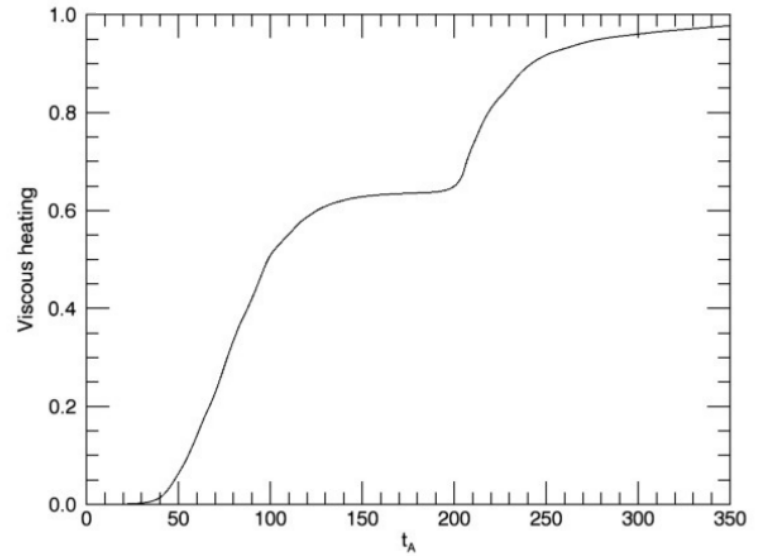
Kinetic energy



Ohmic heating

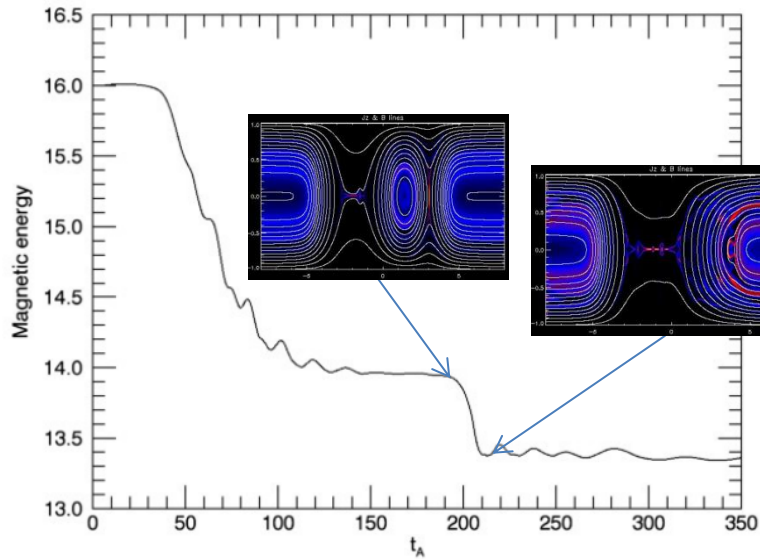


Viscous heating

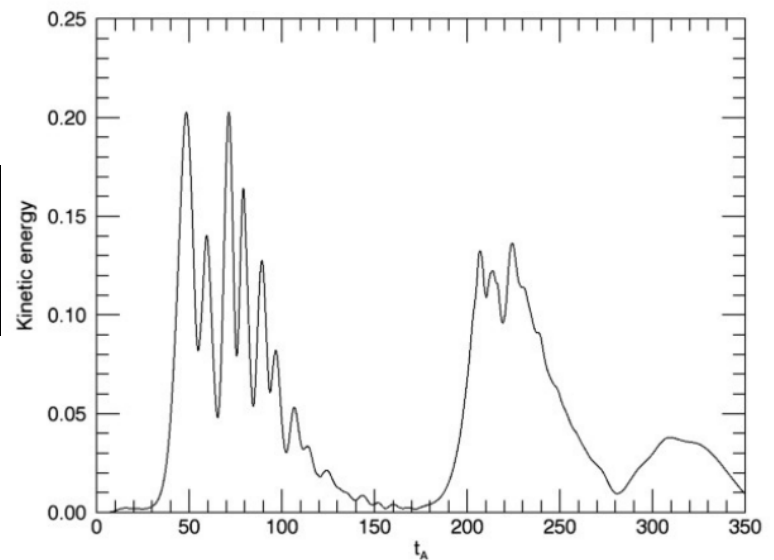


Energetics with merging islands

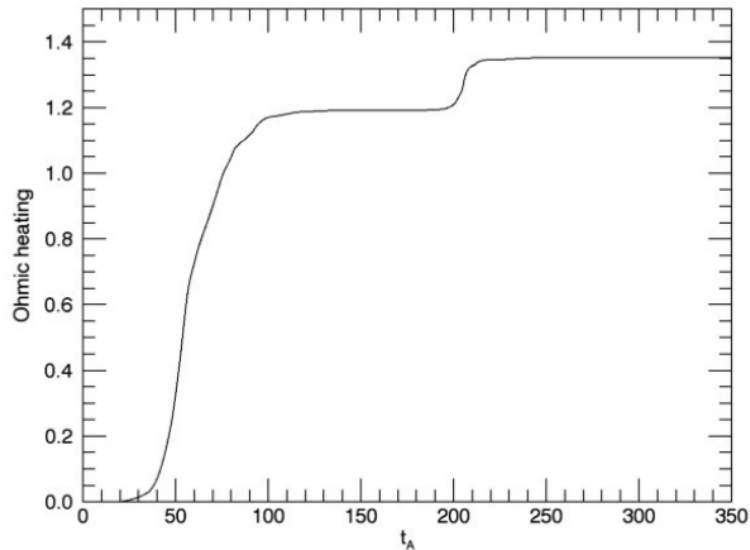
Magnetic energy



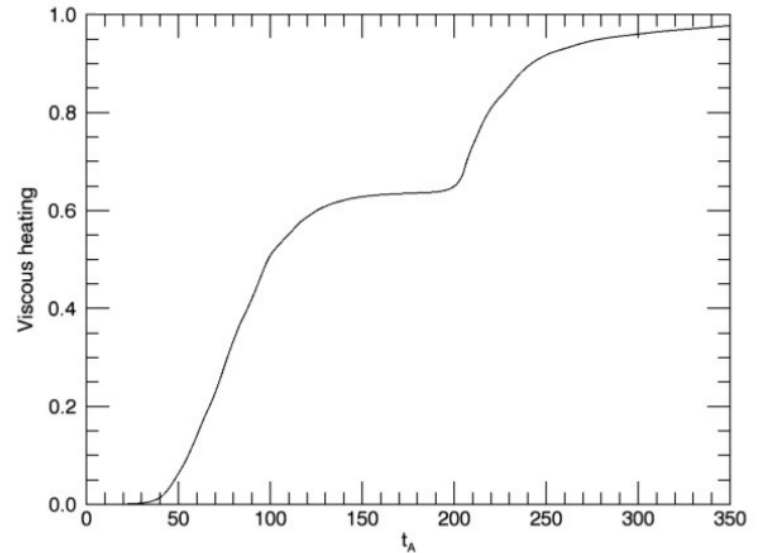
Kinetic energy



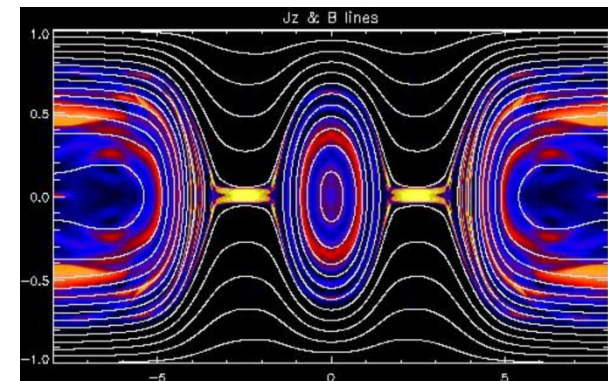
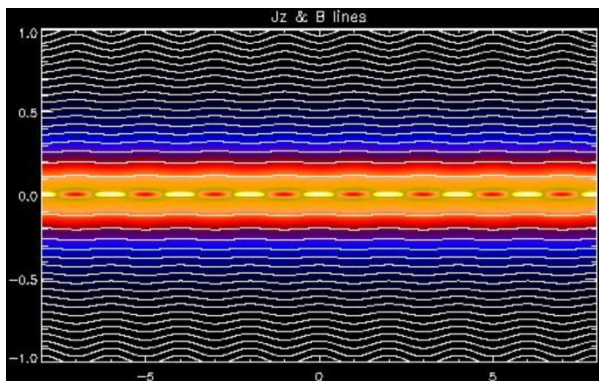
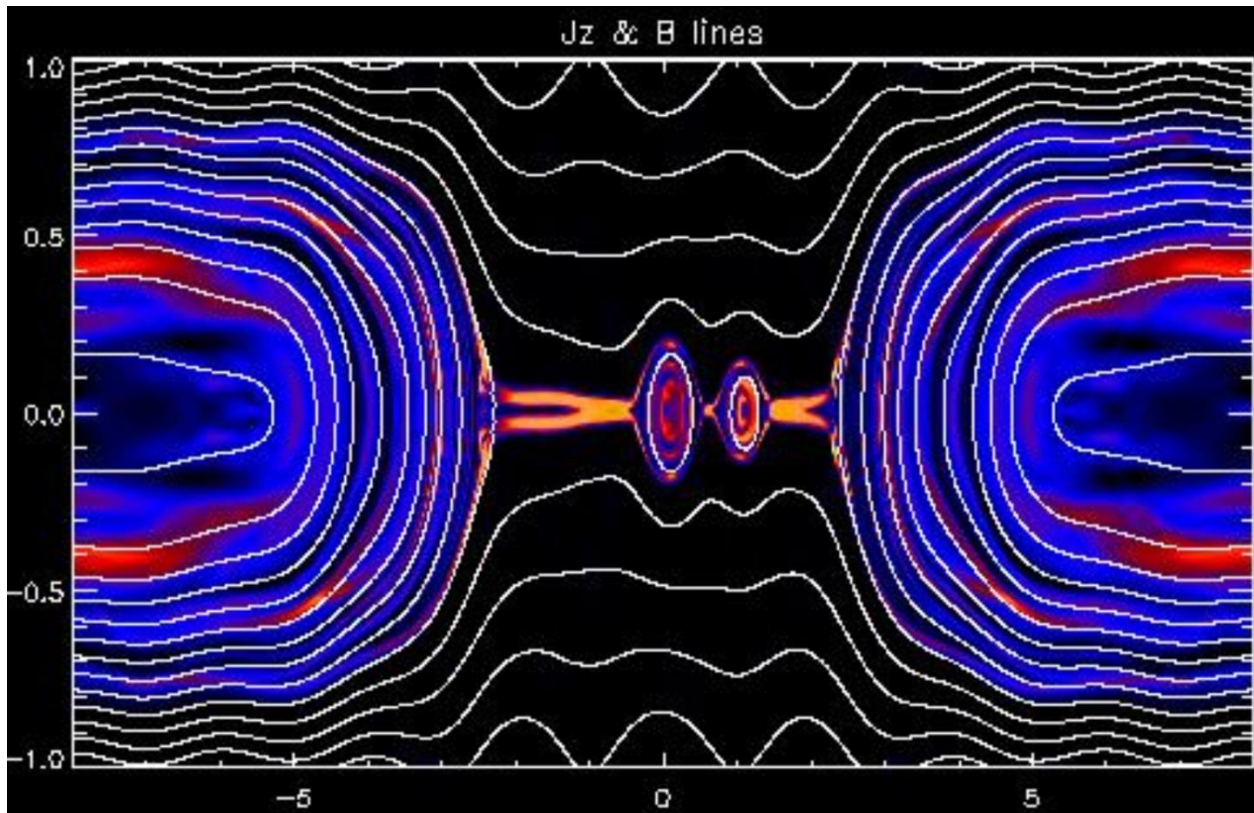
Ohmic heating



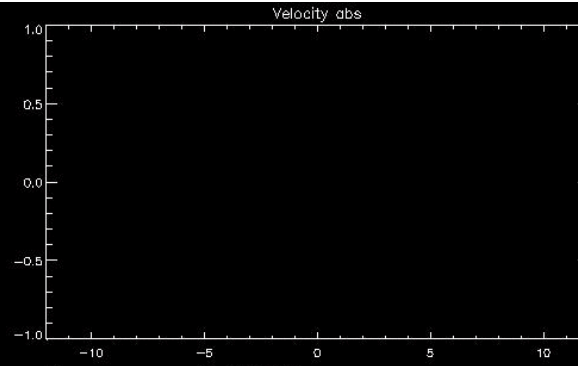
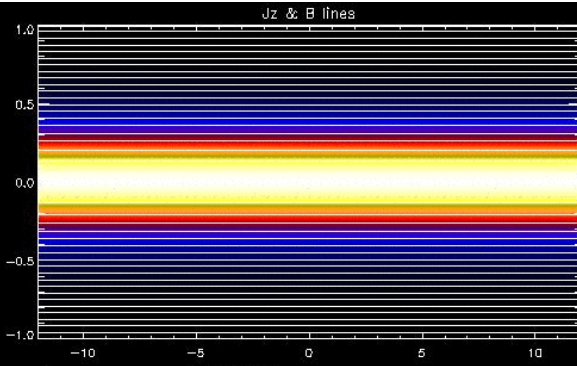
Viscous heating



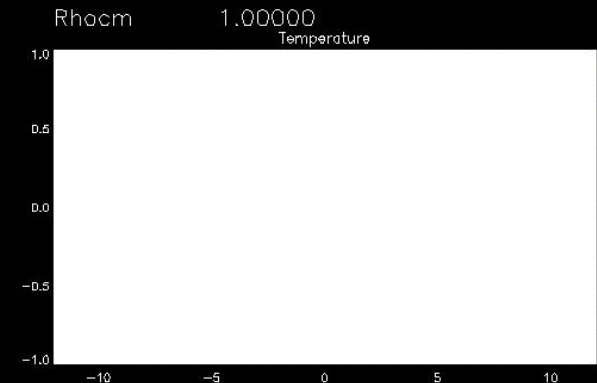
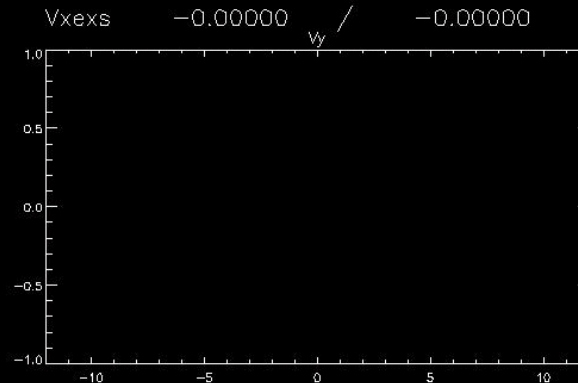
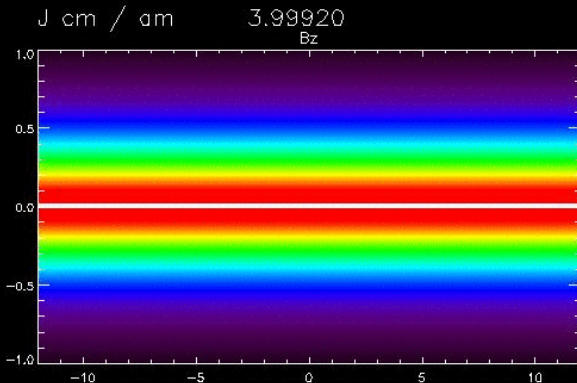
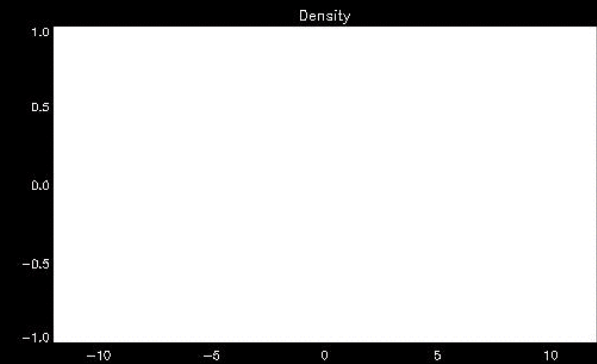
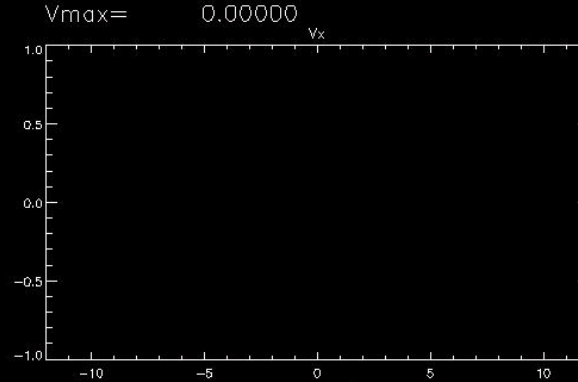
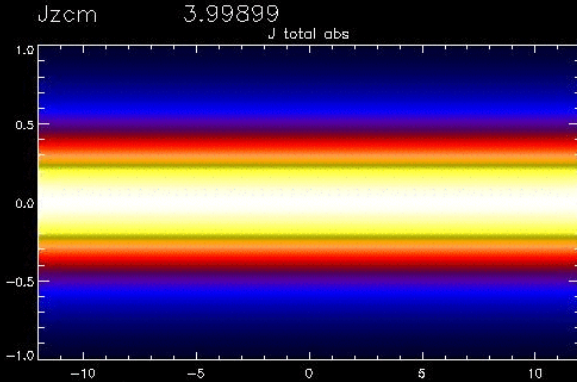
Mixed wavelength perturbation



Mixed wavelength perturbation



Experiment: mp16/005/Data
Time= 0.000000
Max vel= 0.00000
Max current= 3.99920
Max density= 1.00000
Max temp= 0.00250000



J_{cm} 3.99899

V_{max} 0.00000

ρ_{cm} 1.00000

J_{cm} / am 3.99920

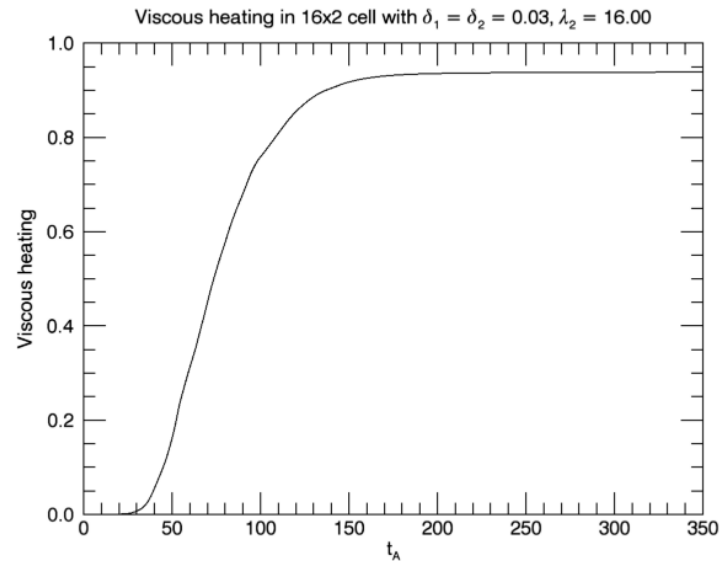
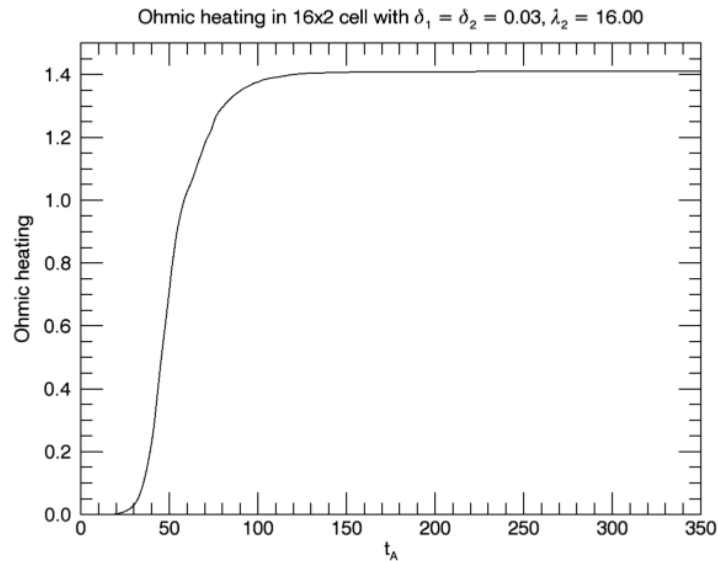
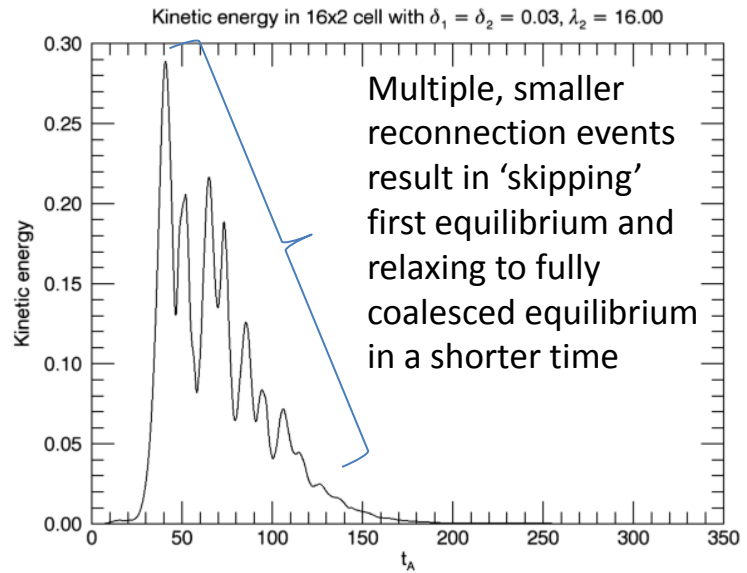
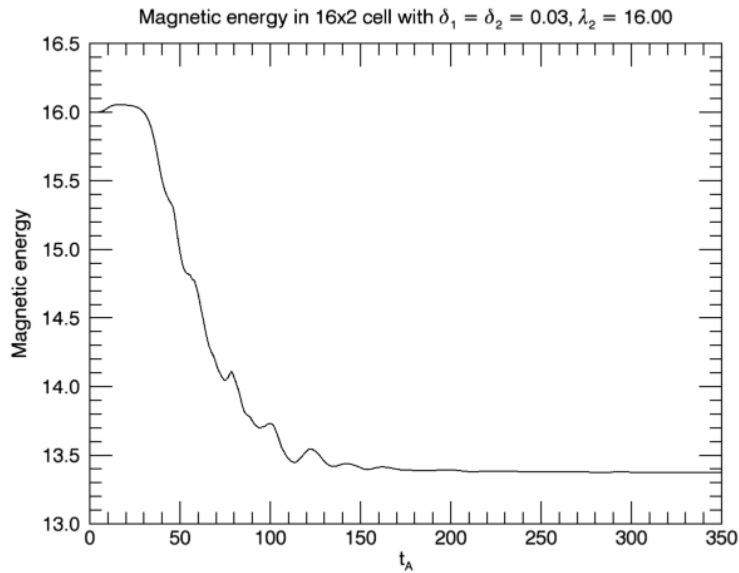
v_{xexs} -0.00000 / -0.00000

T_{cm} 0.00250000

B_{zexs} 0.999946

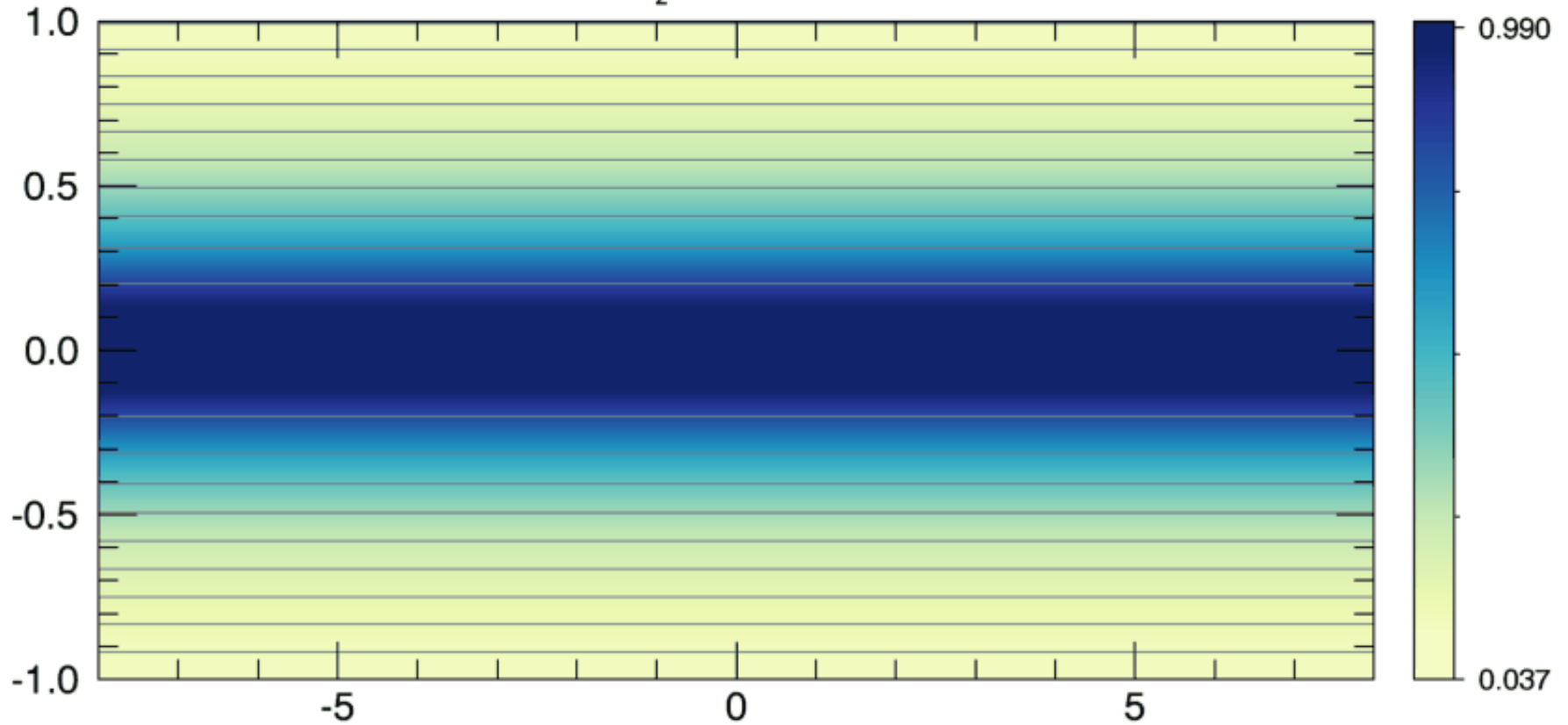
v_{yexs} -0.00000 / -0.00000

Energetics for multiple wavelength perturbation: $L_1 = 16, L_2 = 2$

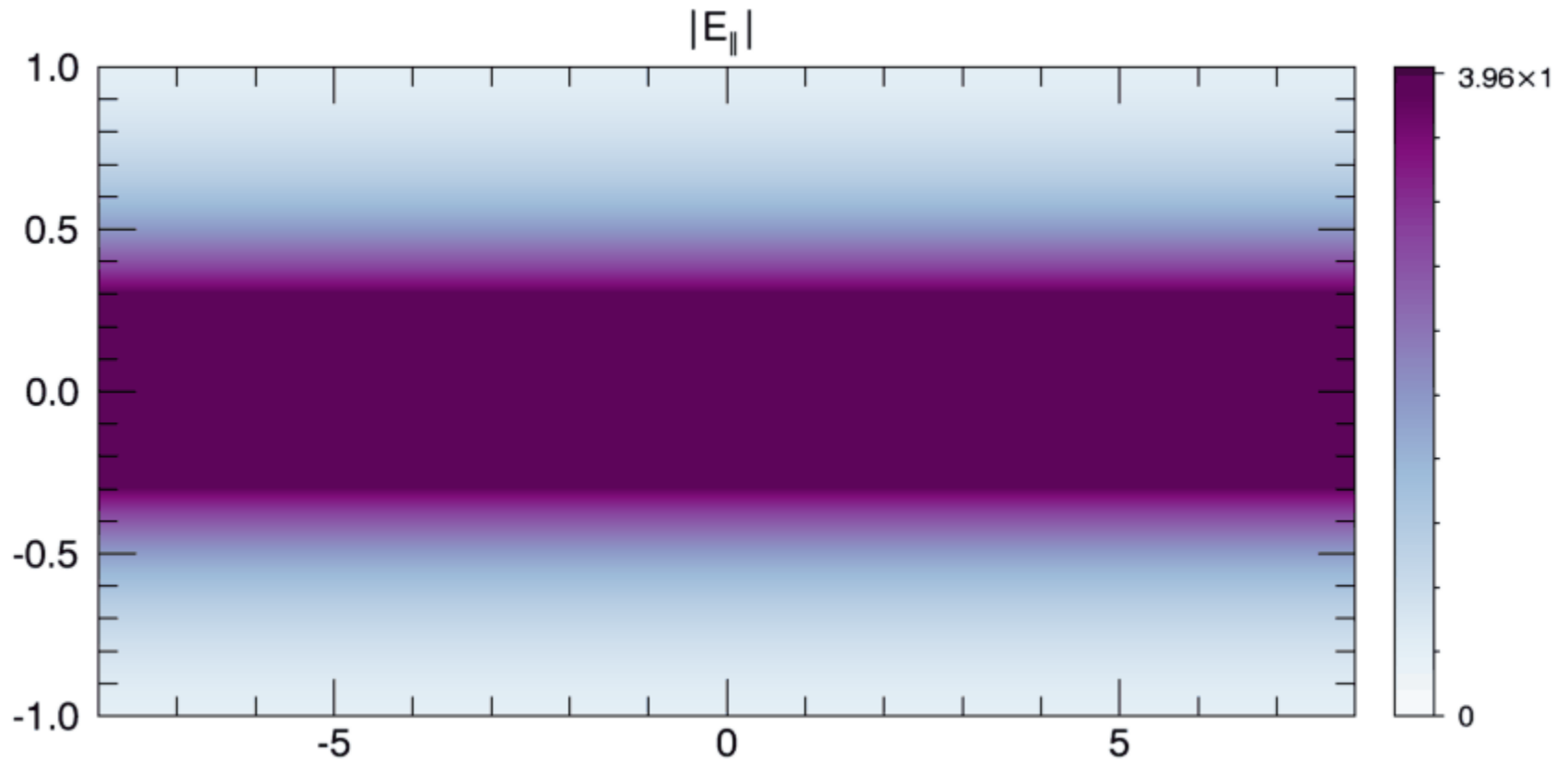


Localised driving perturbation (Gaussian) – Magnetic field

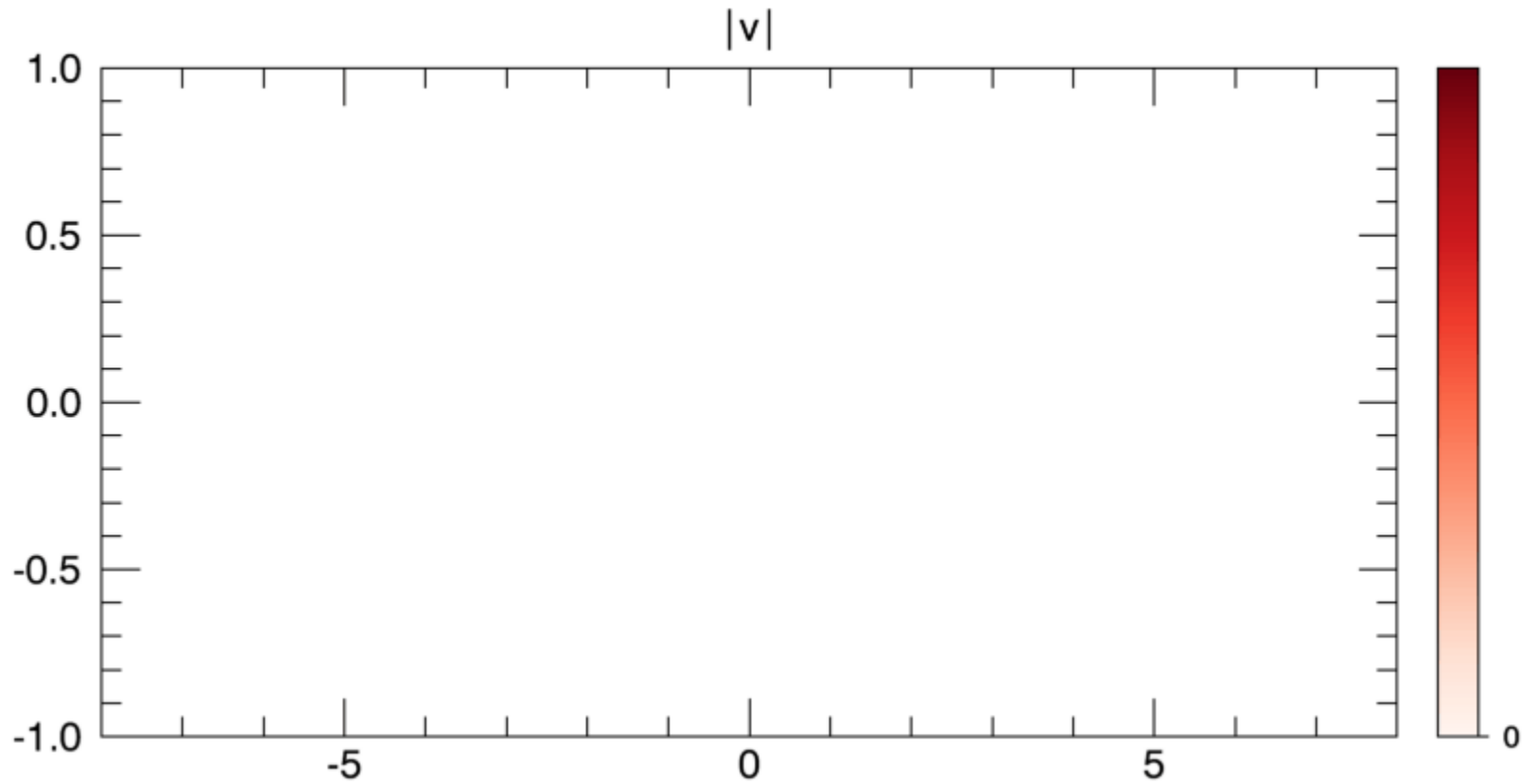
B_z & B lines



Localised driving perturbation (Gaussian) – Electric field



Localised driving perturbation (Gaussian) – velocity field



Test particles

- Test particle trajectories were calculated using GCA (Guiding Centre Approximation) code¹
- Takes time-dependent MHD fields, interpolates fields in time and space.
- Scale parameters:
 $L_0 = 10^4 \text{ m}$ $B_0 = 3 \times 10^{-3} \text{ T}$ $\rho = 2.4 \times 10^9 \text{ cm}^{-3}$
- Boundary conditions: periodic in x, free in y & z.

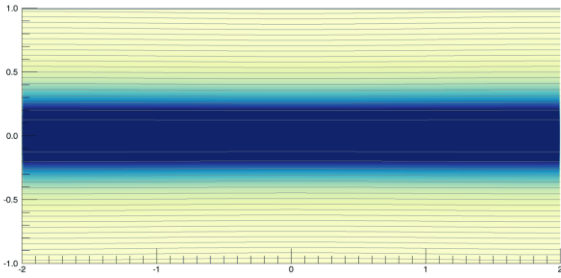
¹ Gordovskyy et al., ApJ 2010

² Northrop, 1963

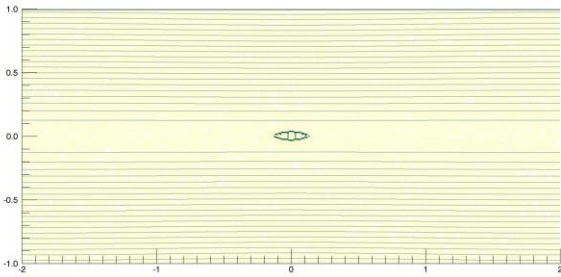
Short domain trajectories

No coalescence

$|J_z|$ & B lines



$|E_{||}|$ & B lines

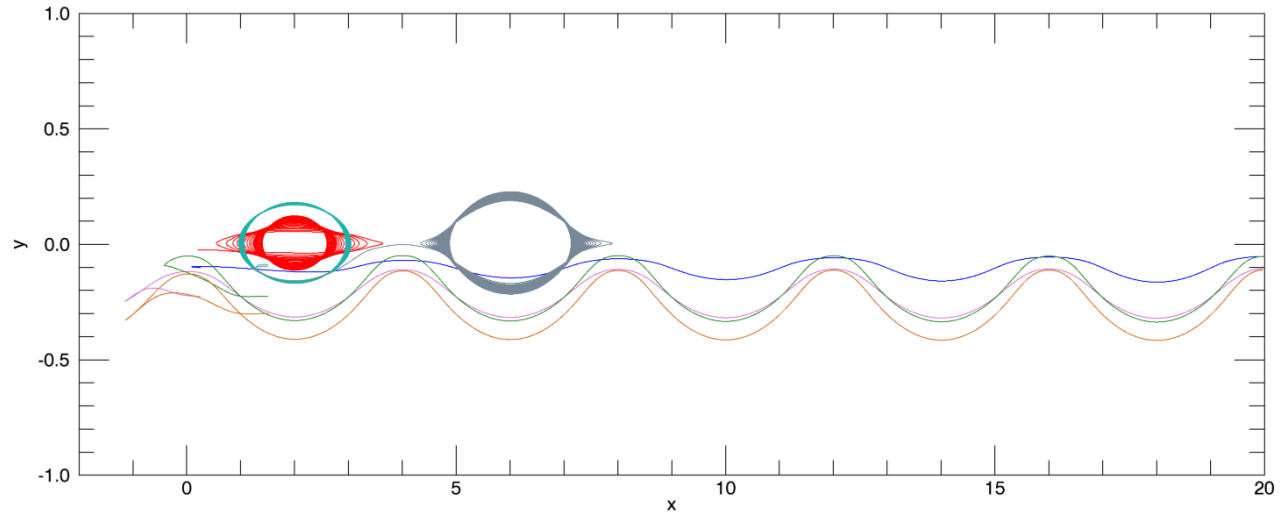


MHD parameters:

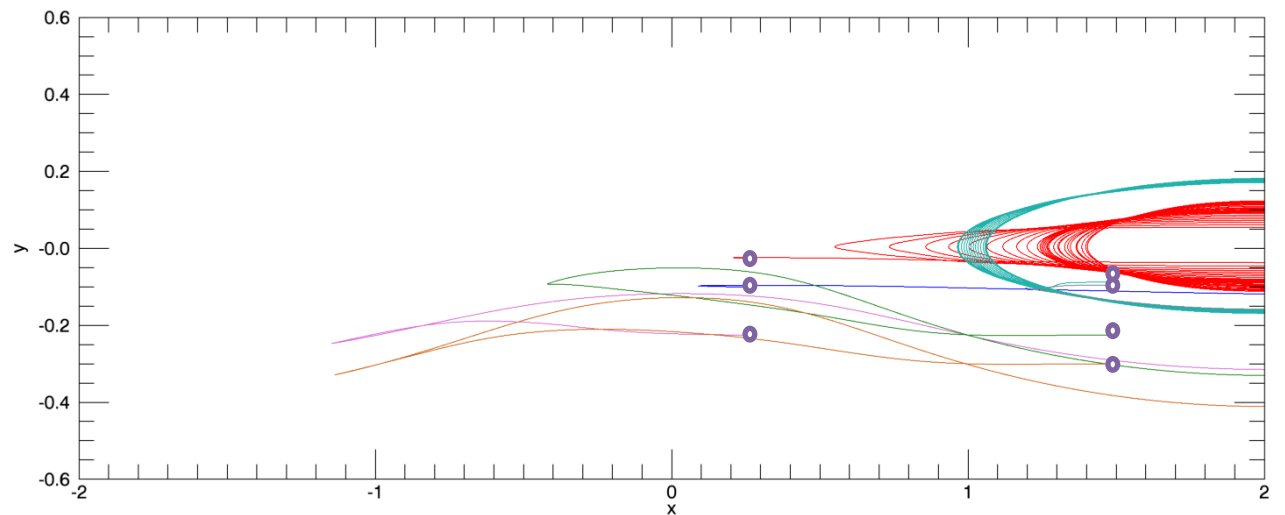
- Initial current sheet thickness $y_0 = 0.25$
- $\eta = 3.2 \times 10^{-4}$ when $J > 4.1$, else $\eta = 0$

Right: Bottom plot is a zoomed in section of top. Initial particle positions highlighted by dots.

Typical trajectories: **x-y** plane

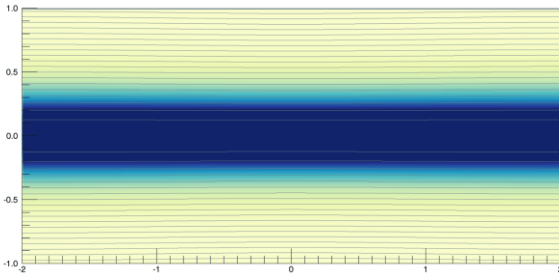


Typical trajectories: **x-z** plane

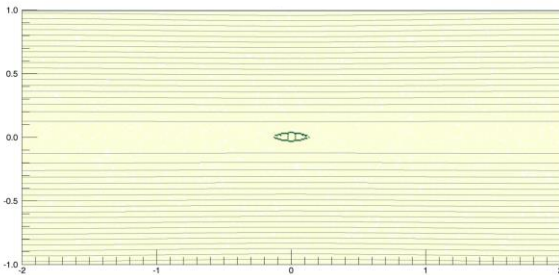


Short domain trajectories

$|J_z|$ & B lines



$|E_{//}|$ & B lines

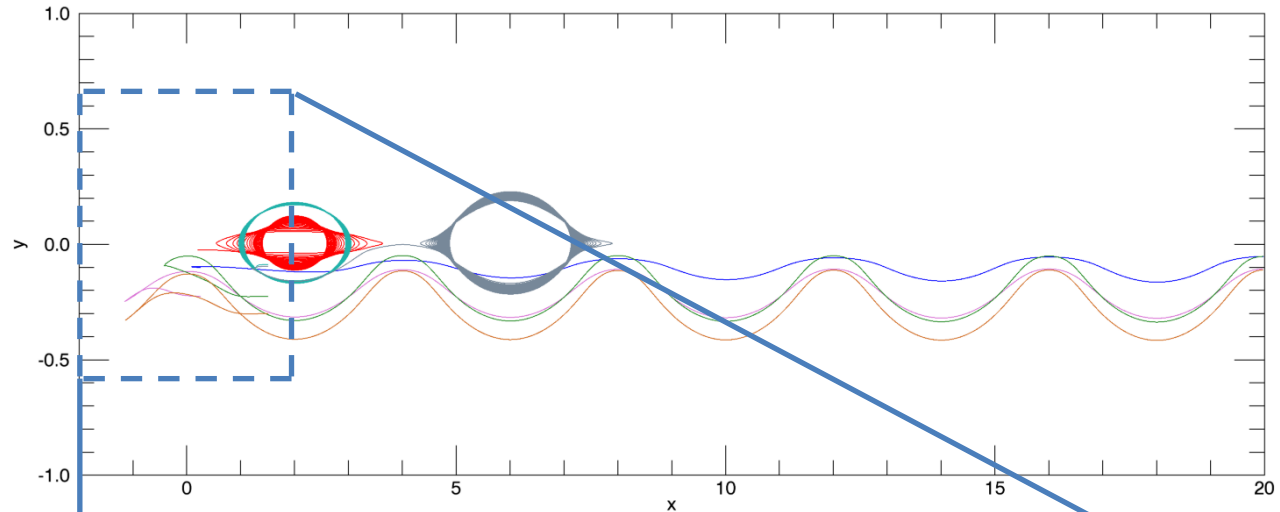


MHD parameters:

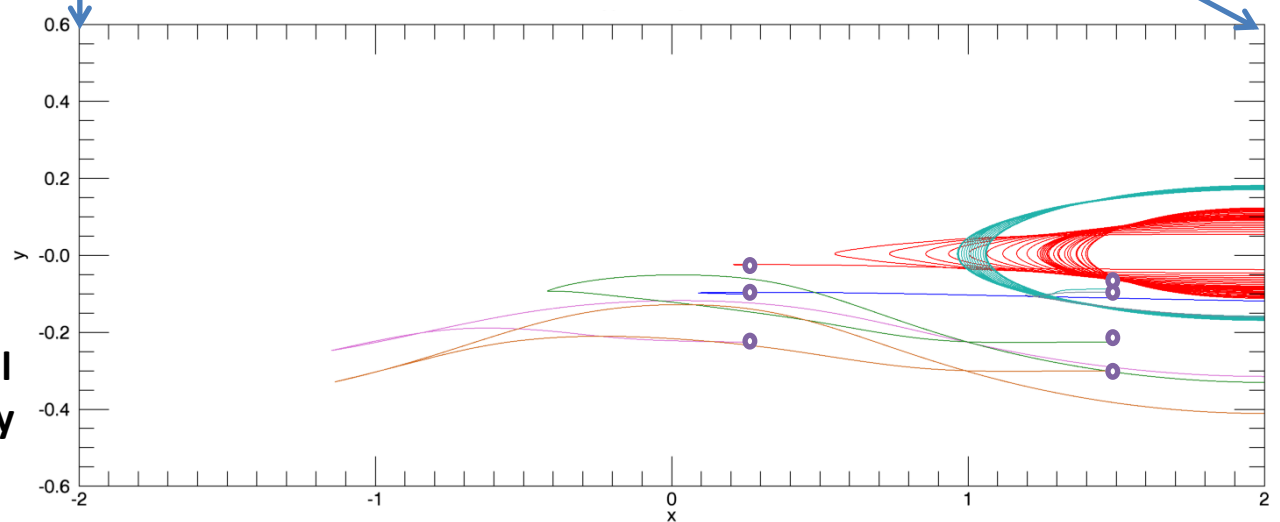
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Zoomed in section of top. Initial particle positions highlighted by dots.

Typical trajectories: x - y plane



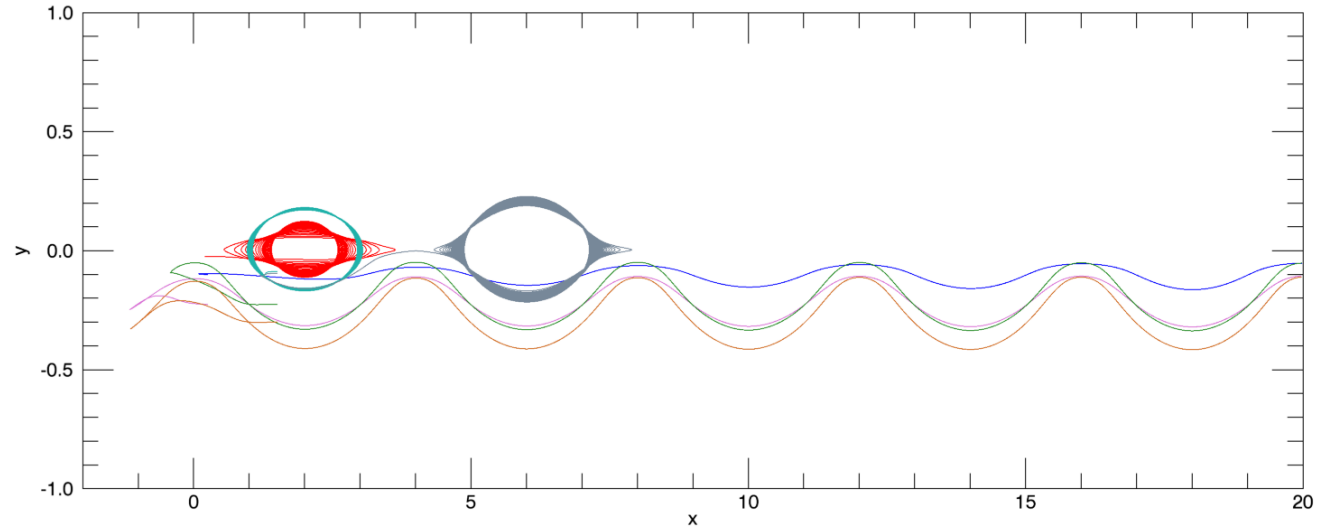
Typical trajectories: x - y plane



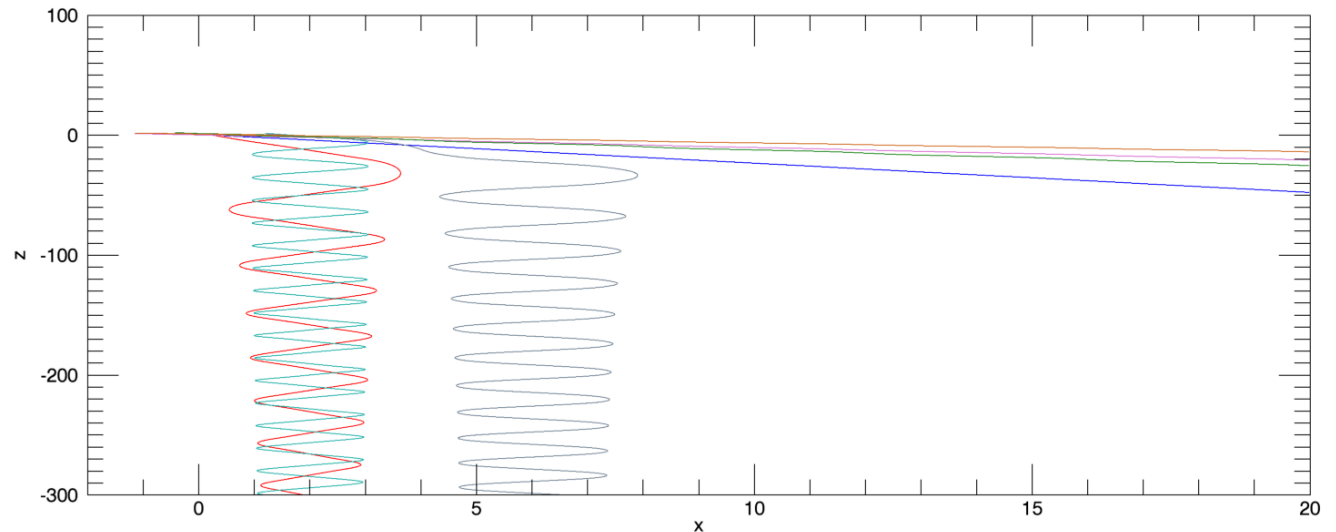
Short domain trajectories

Particles on closed field lines travel further in z direction, indicating greater acceleration due to concentrated electric field at island separatrices.

Typical trajectories: x - y plane



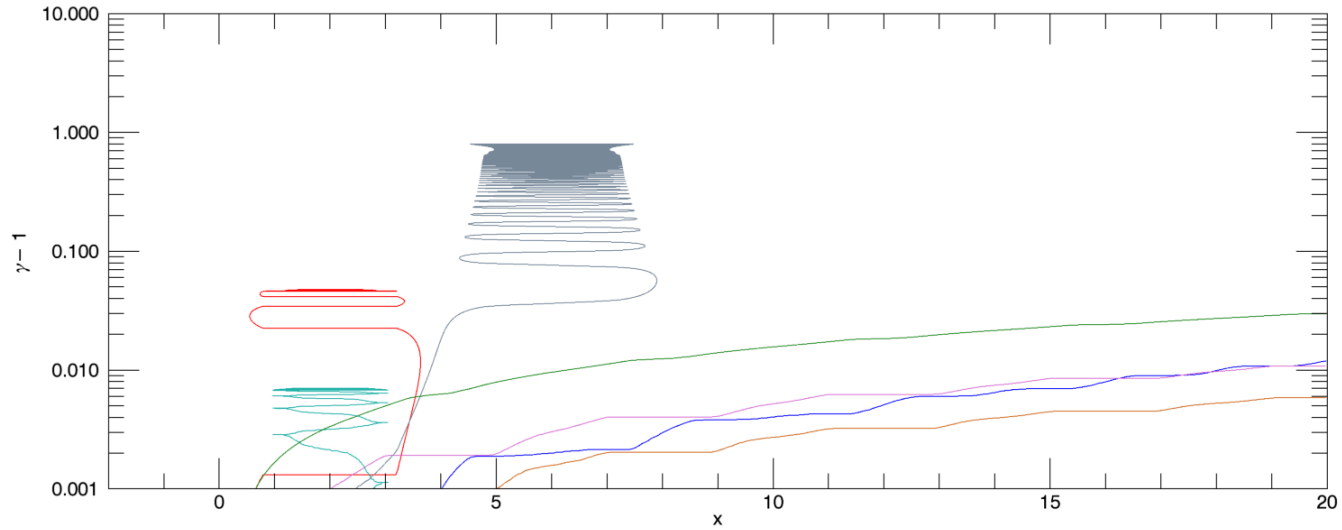
Typical trajectories: x - z plane



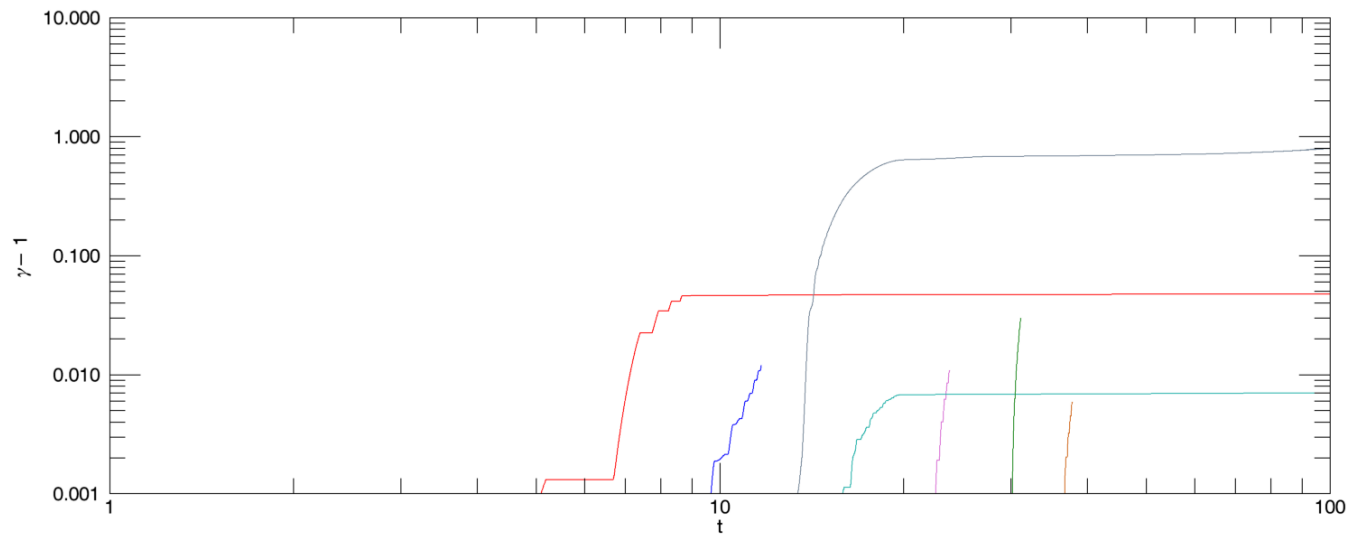
x - z plane

Short domain trajectories

Particle energy against x position

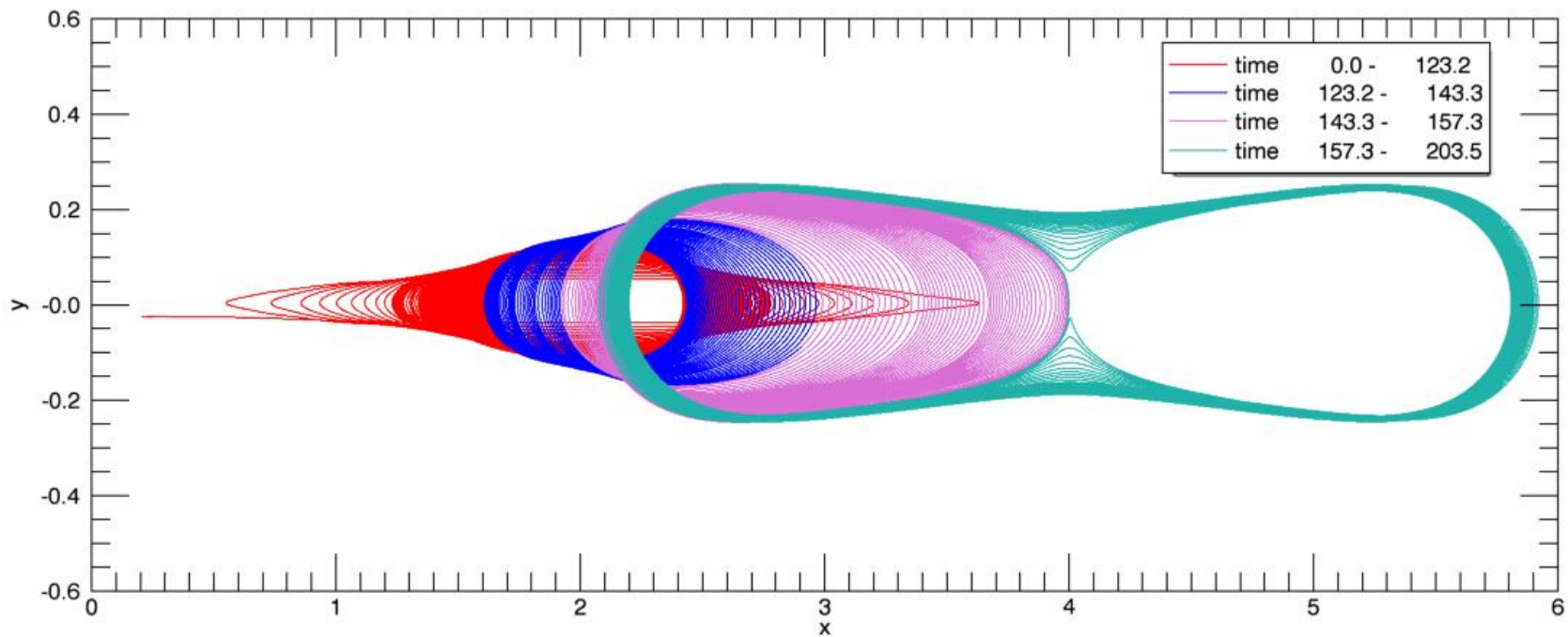


Particle energy against x position

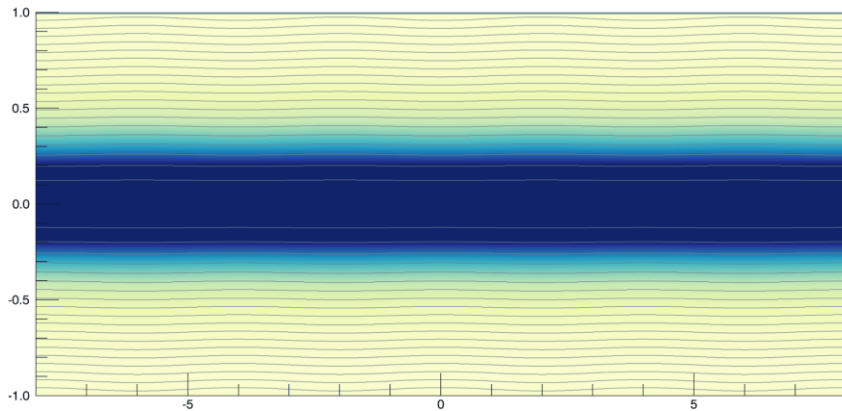


Long domain with island coalescence - trajectories

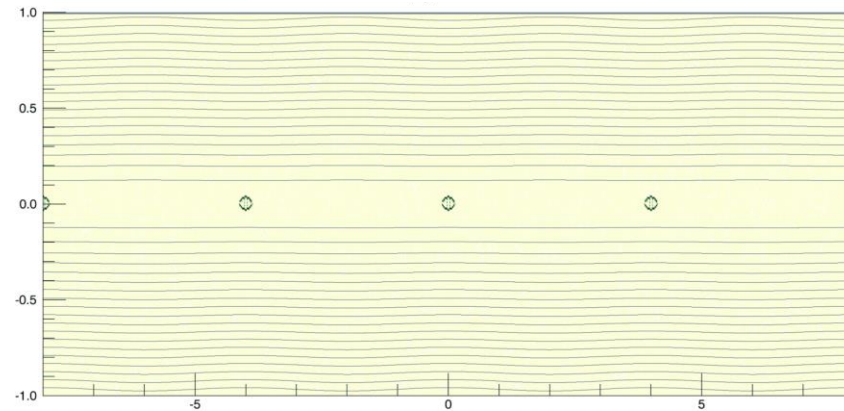
Position in x - y plane



$|J_z|$ & B lines

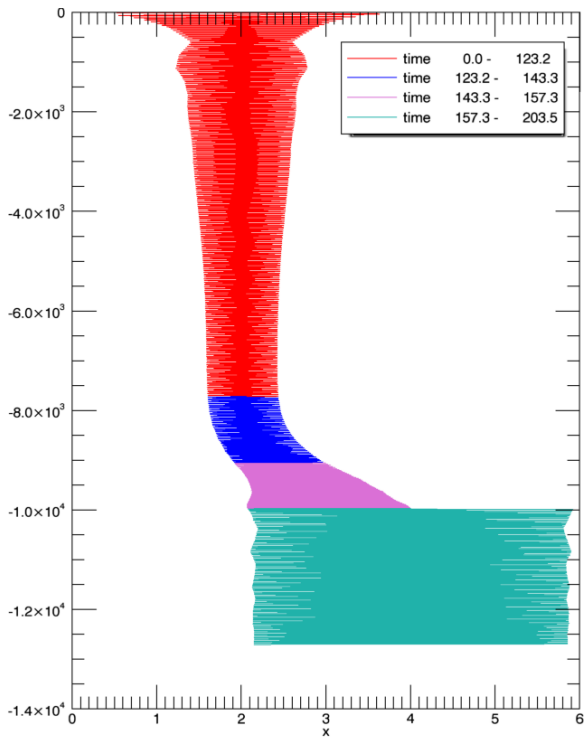


$|E_{||}|$ & B lines

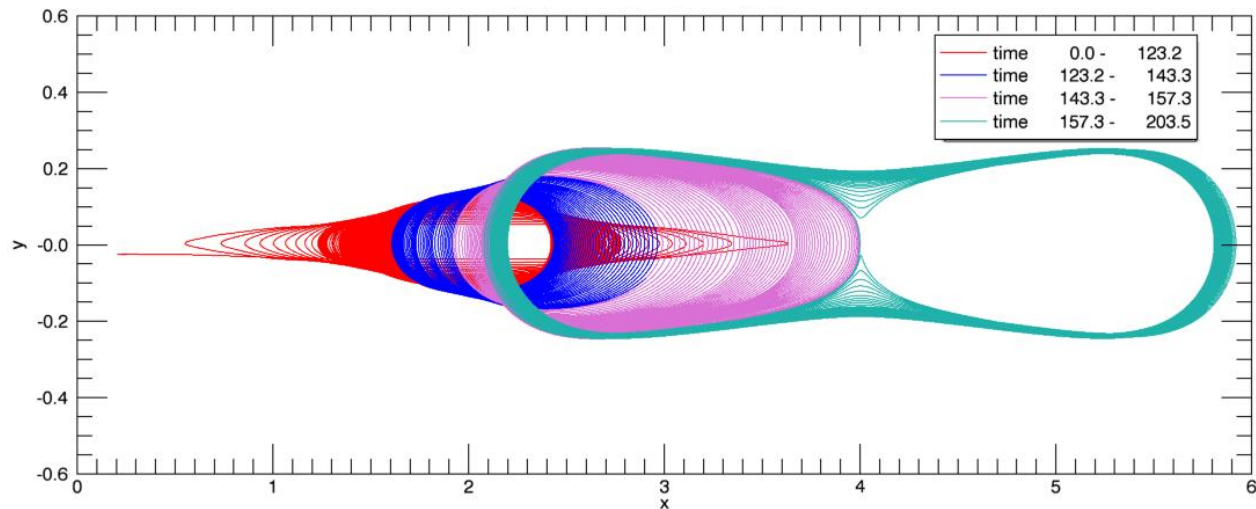


Long domain with island coalescence - trajectories

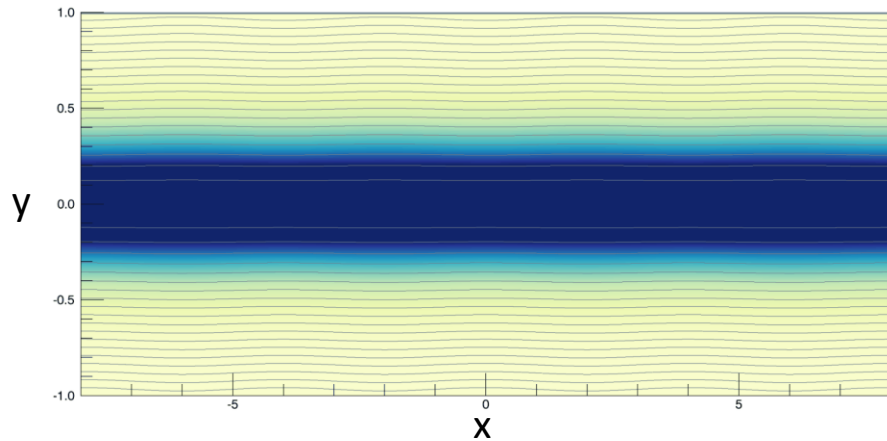
Position in x - z plane



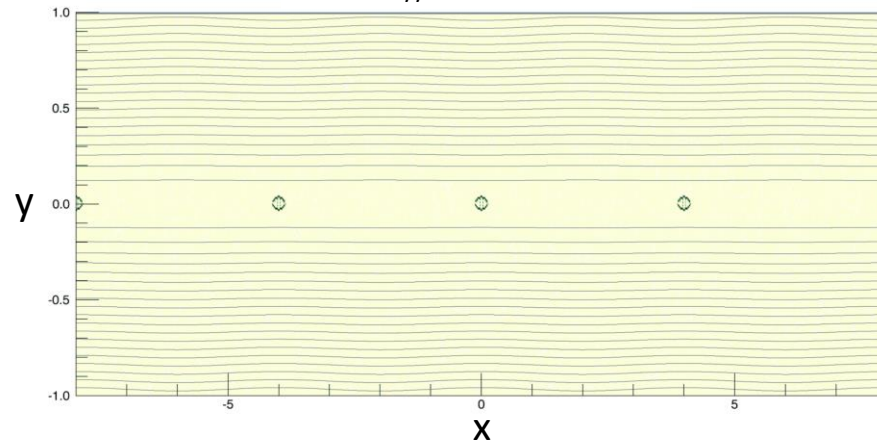
Position in x - y plane



$|J_z|$ & B lines

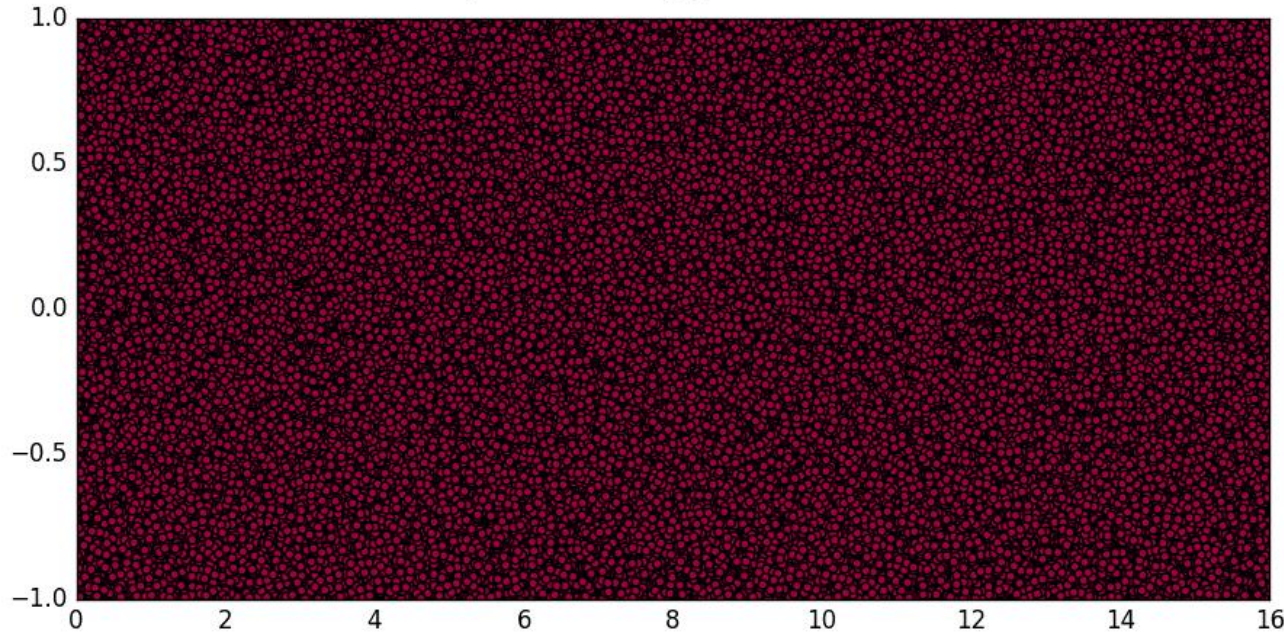


$|E_{//}|$ & B lines

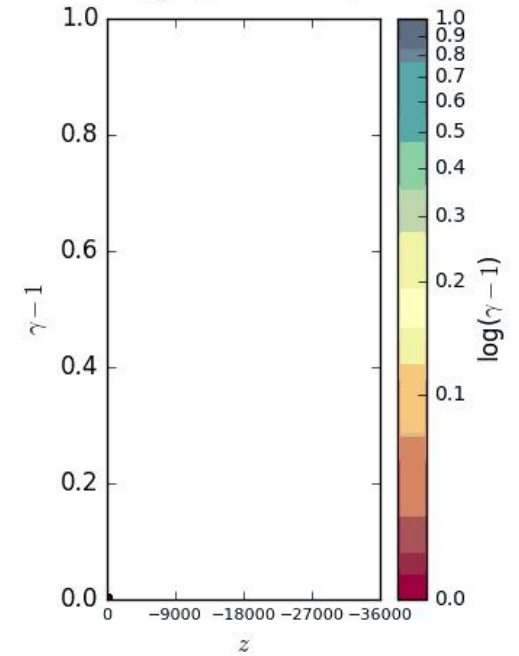


Long domain with island coalescence –ion distribution

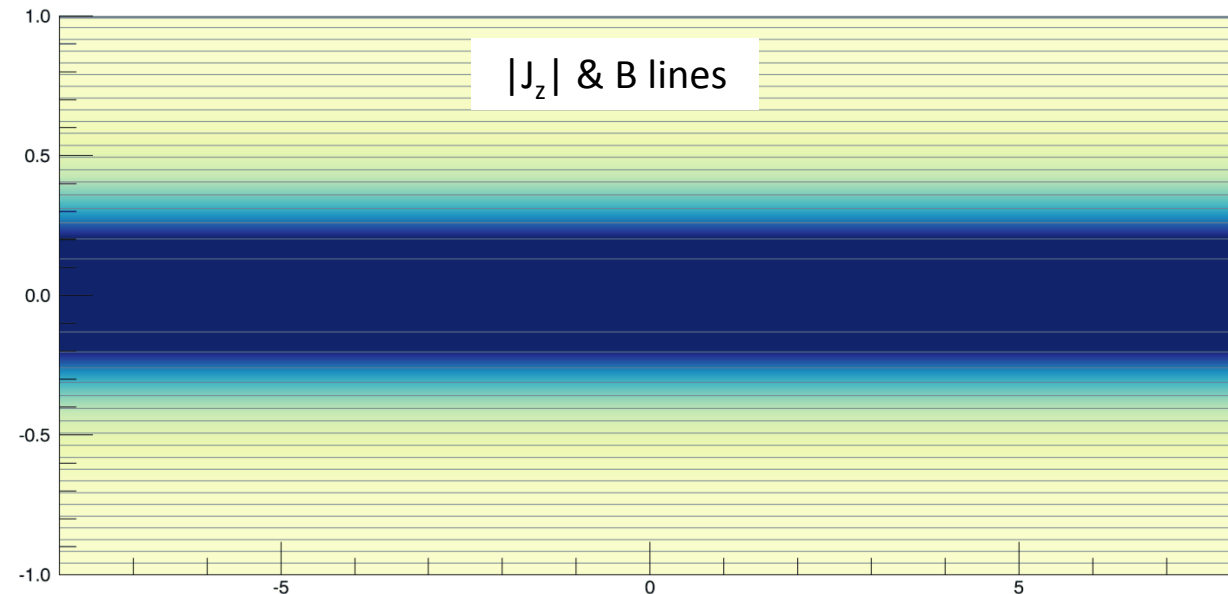
Final positions in x-y plane, $t = 0.0$



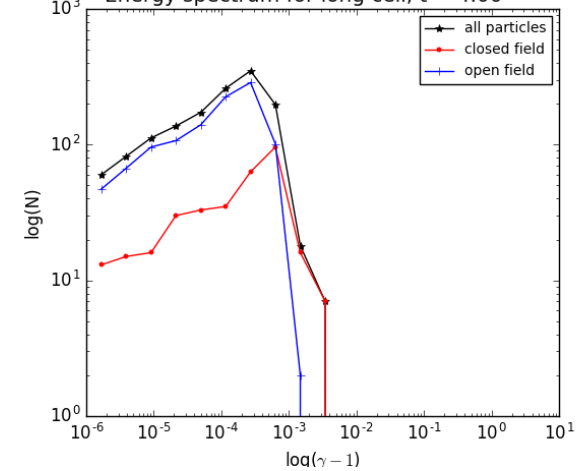
Energy against height z



$|J_z|$ & B lines



Energy spectrum for long cell, $t = 4.00$



Summary & future work

- Particles on closed field (magnetic islands) are preferentially accelerated due to concentration of electric field about separatrices
- Island coalescence second step of magnetic reconnection & particle acceleration.
- Electric fields in coalescence are very localised and in reverse direction
- Forced reconnection and island coalescence could play a role in multi-stage acceleration in solar flares

- Further analysis of particle energisation in 2D fields
- 3D geometries
 - Islands \rightarrow twisted flux ropes (kink instability?)
 - Reconnection on different resonant surfaces \rightarrow island overlap (stochastic fields?)